

Bispectrum of the Sunyaev-Zeldovich Effect

Suman Bhattacharya

Kavli Institute for Cosmological Physics
The University of Chicago

&

High Energy Physics Division
Argonne National Laboratory

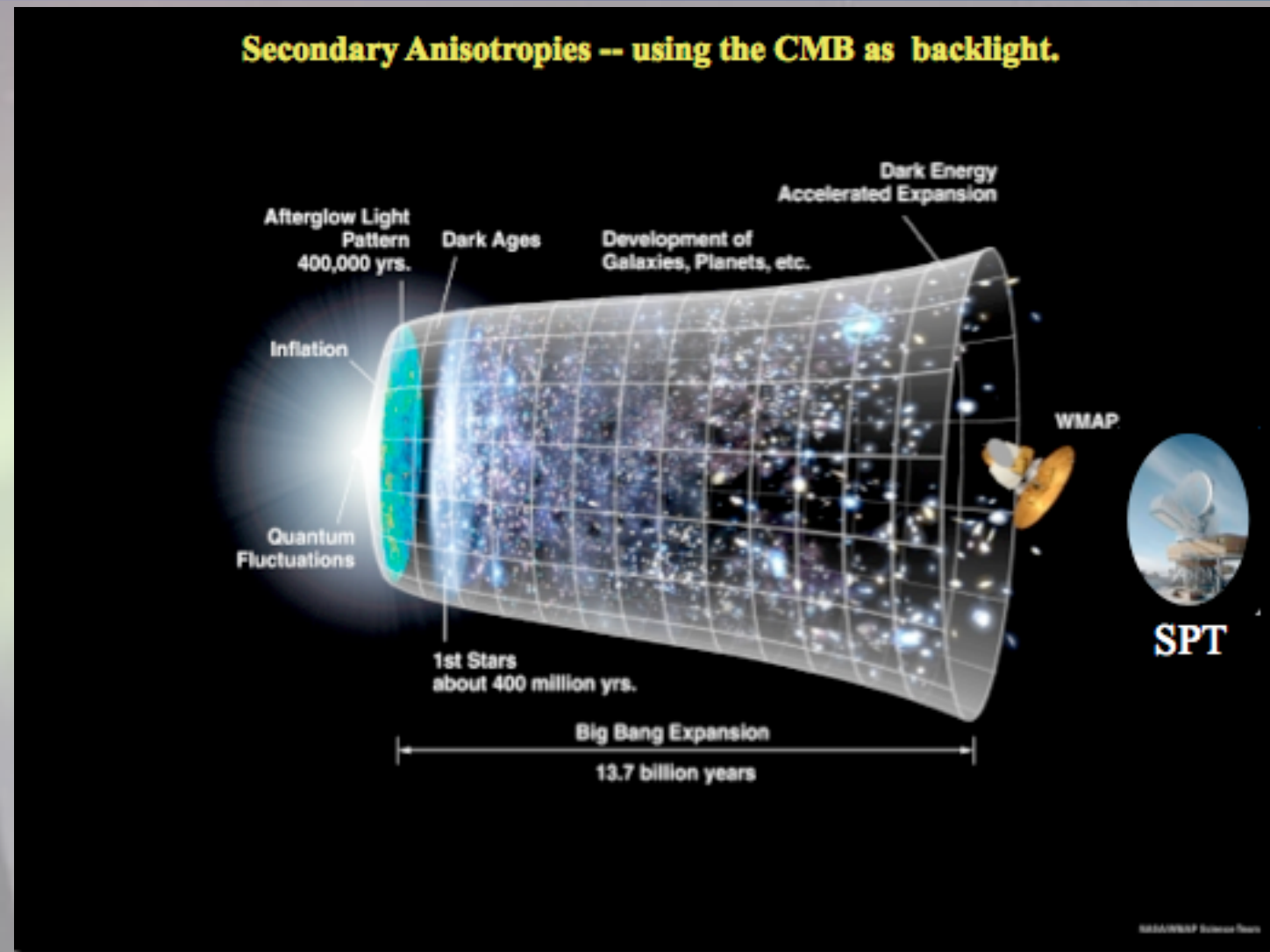
with

Daisuke Nagai, Laurie Shaw (Yale), Gil Holder
(McGill), Tom Crawford (U Chicago)
and SPT Team

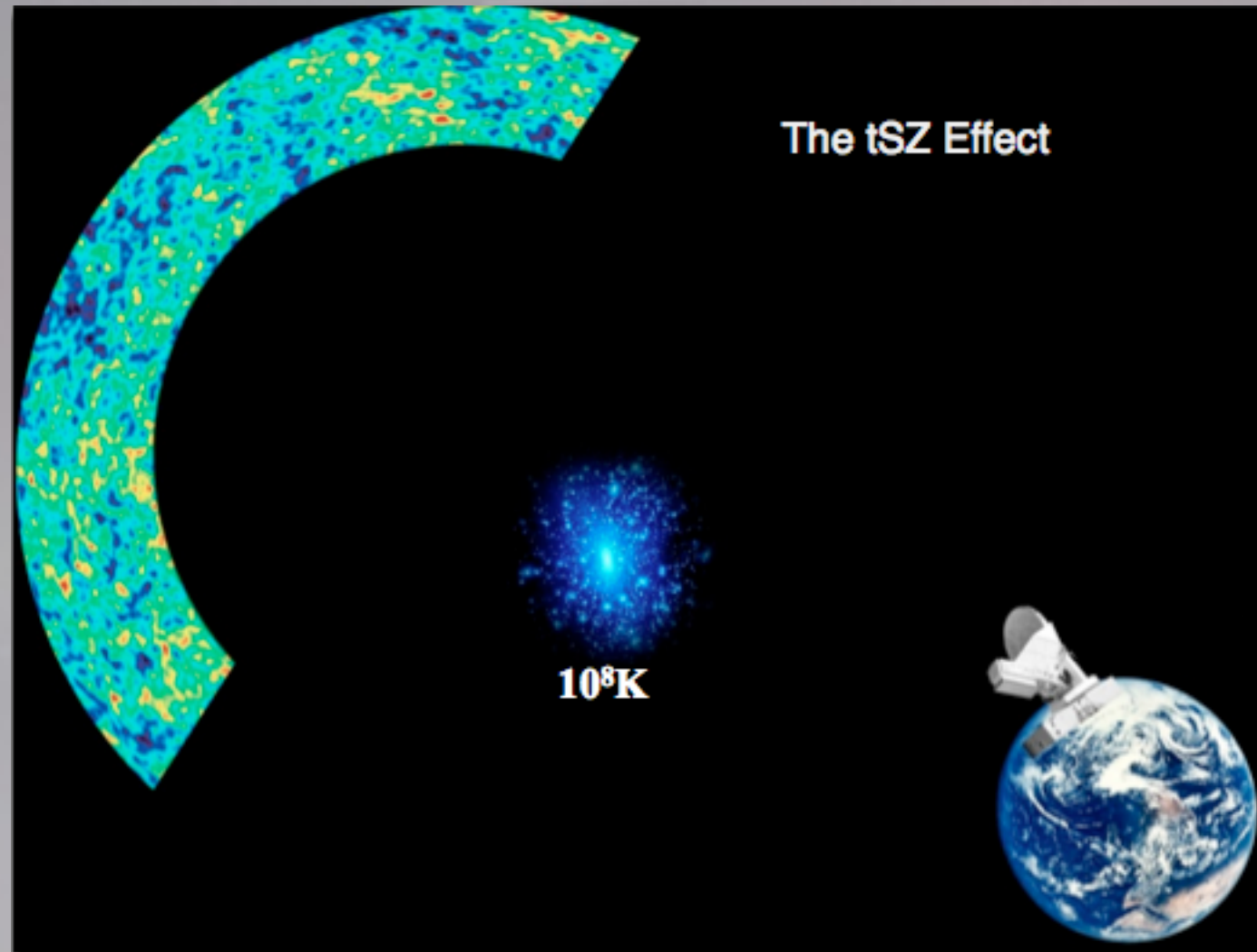
Cosmology meeting, Santa Fe, July, 2013

CMB Science- Secondary Anisotropies

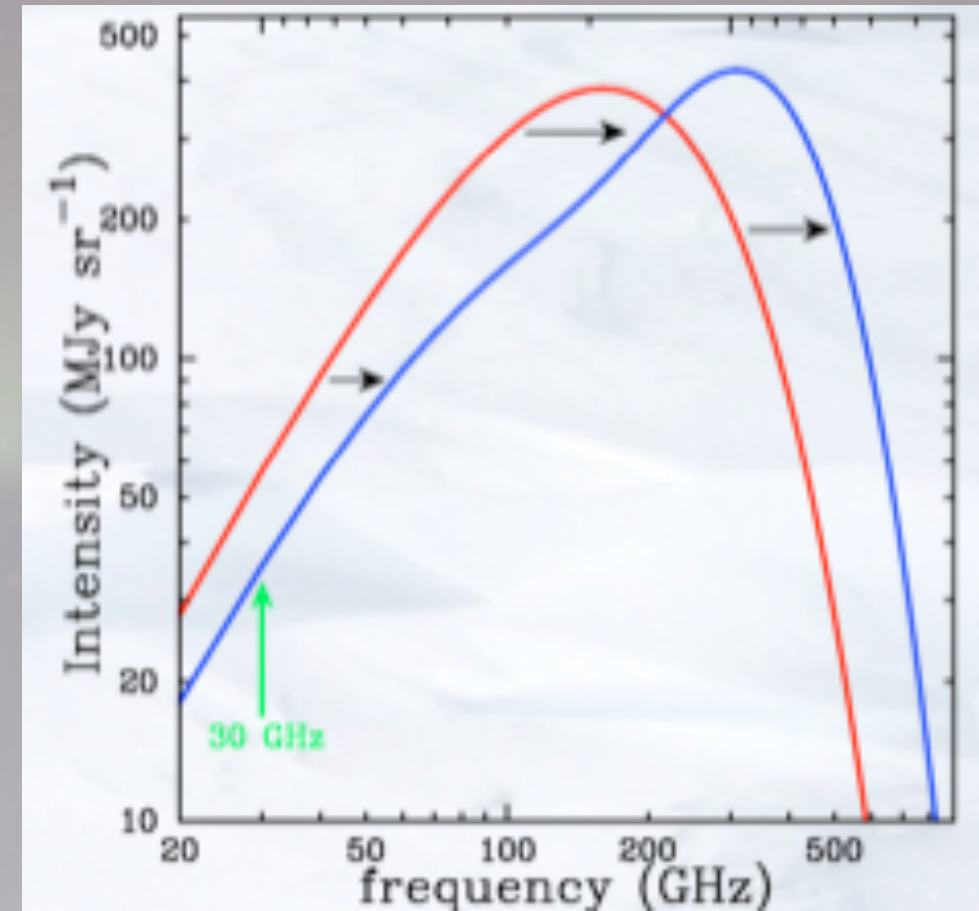
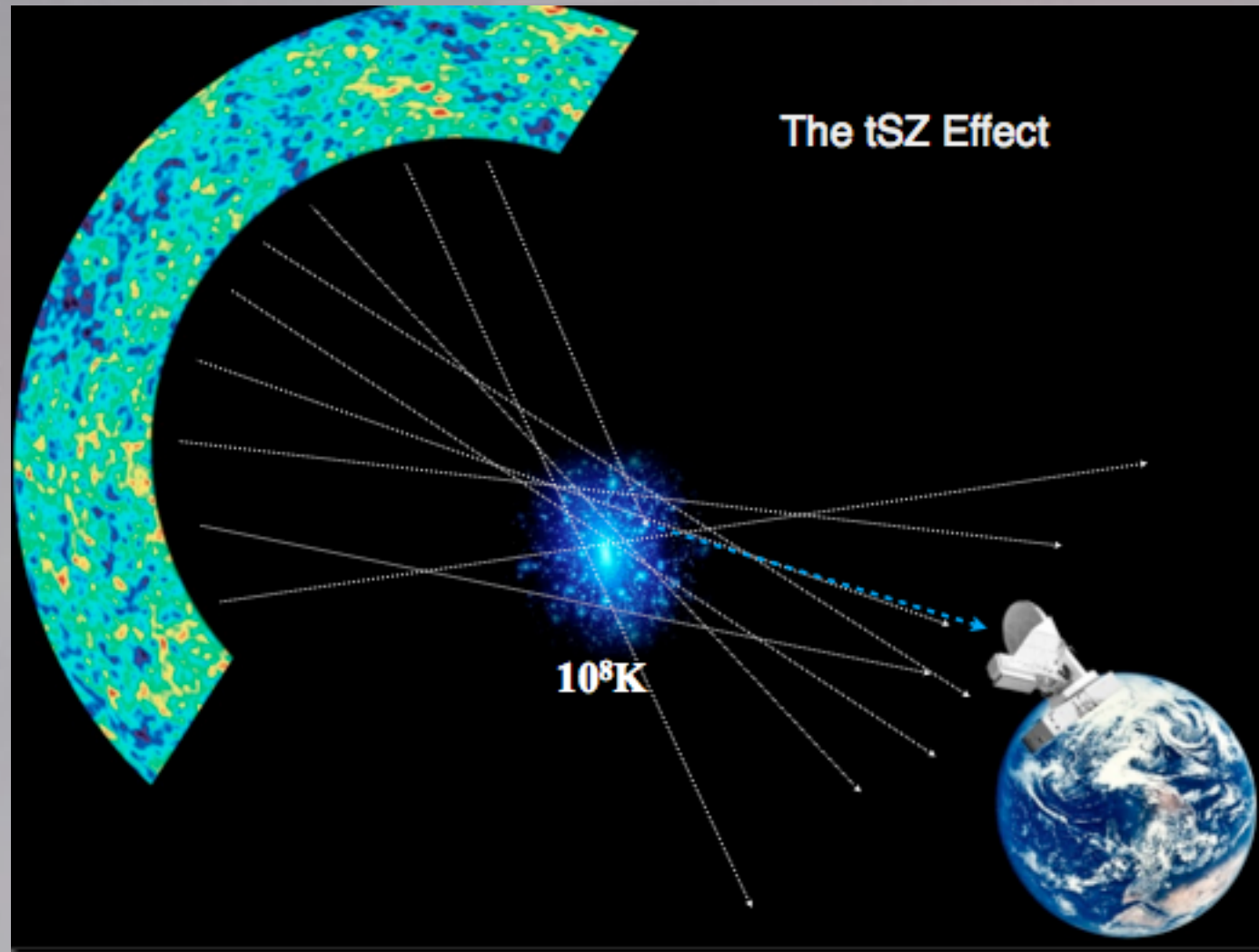
- **Secondary Anisotropies** => late universe, small scale phenomena
- **Sunyaev-Zeldovich Effect**
 - Thermal SZ
 - kinetic SZ
- **CMB Lensing**
- **point sources**
 - radio (AGN), dusty high-z star forming galaxy
 - Poisson , clustered



Secondary Anisotropies-Sunyaev-Zeldovich Effect



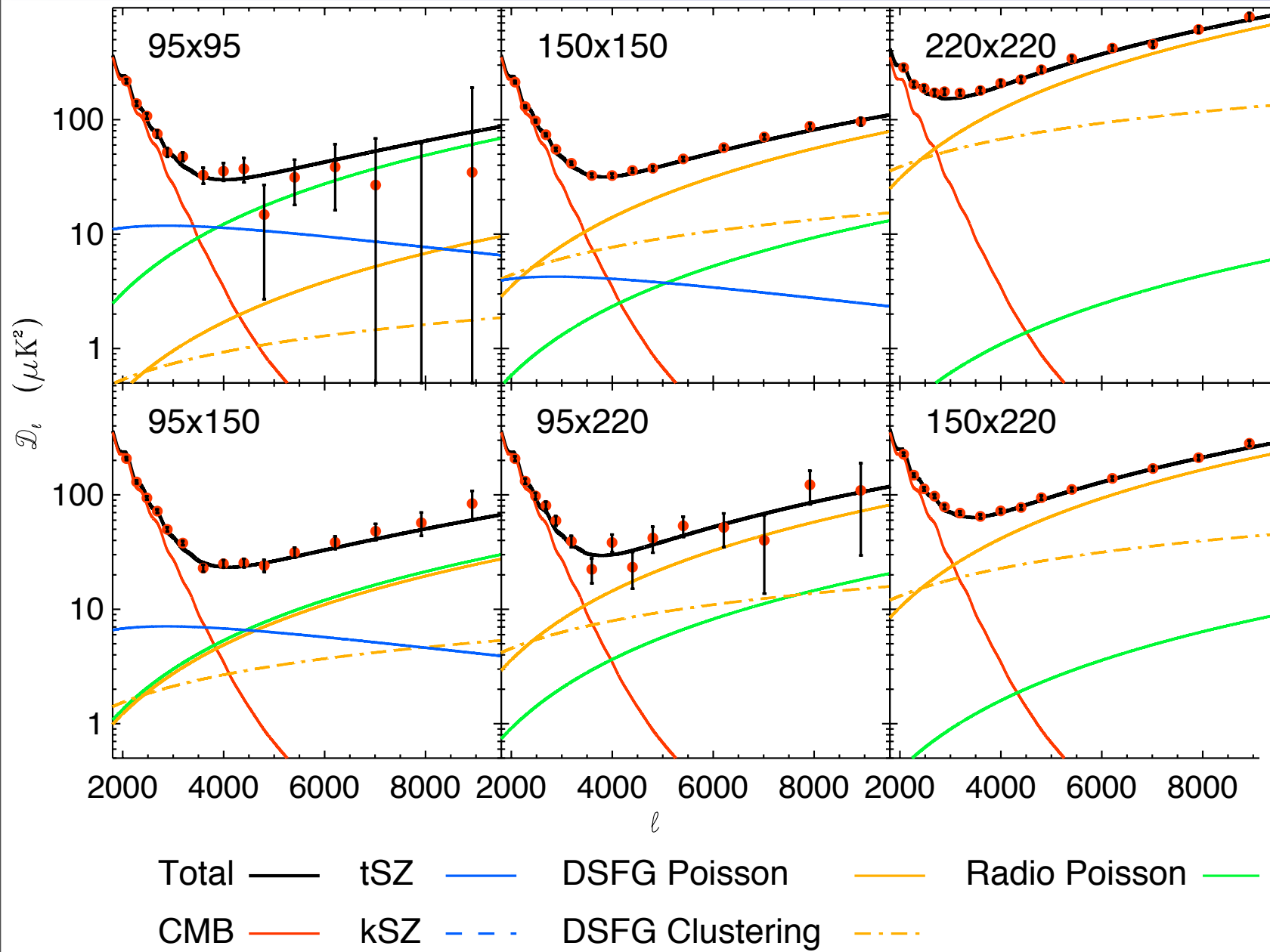
Secondary Anisotropies-Sunyaev-Zeldovich Effect



- 1-2 % CMB photons passing through galaxy clusters get inverse Compton scattered to higher energy
- Surface Brightness is independent of redshift

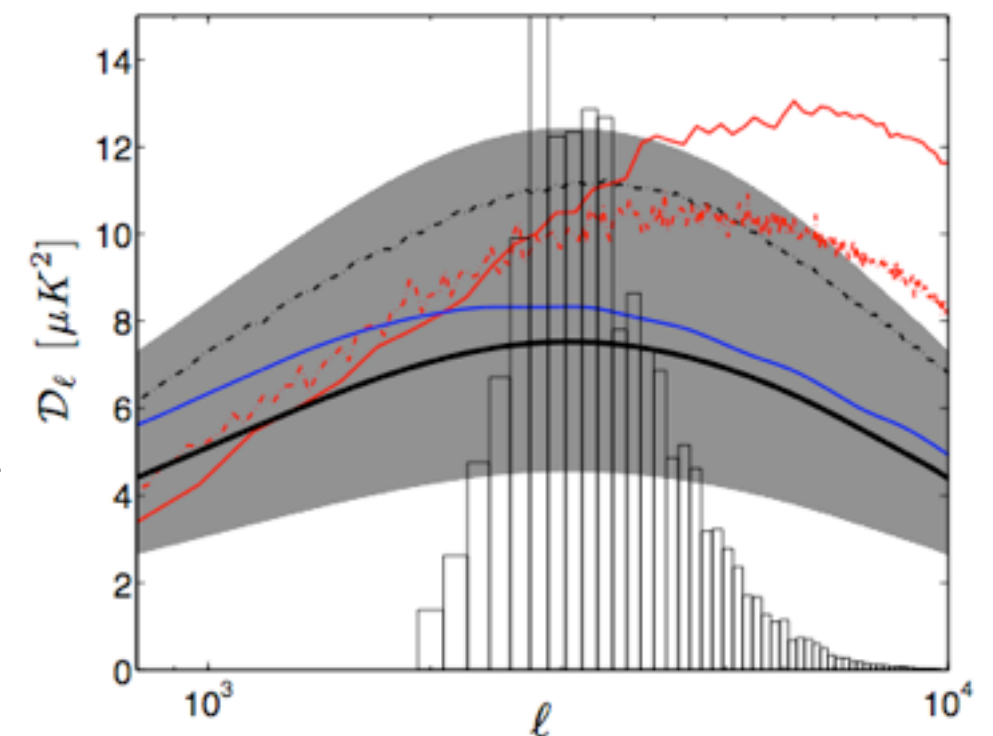
$$\frac{\Delta T_{cmb}}{T_{cmb}} \equiv f_{\nu}(x)y = \left(\frac{k_B \sigma_T}{m_e c^2} \right) \int n_e(l) T_e(l) dl$$

SZ spectrum studies



Reichardt et al, SPT team

Theory uncertainty



Theory uncertainty is the limiting factor!



Pressure Profile: Theoretical Model

Modeling pressure profiles is the key to understand the SZ effect !

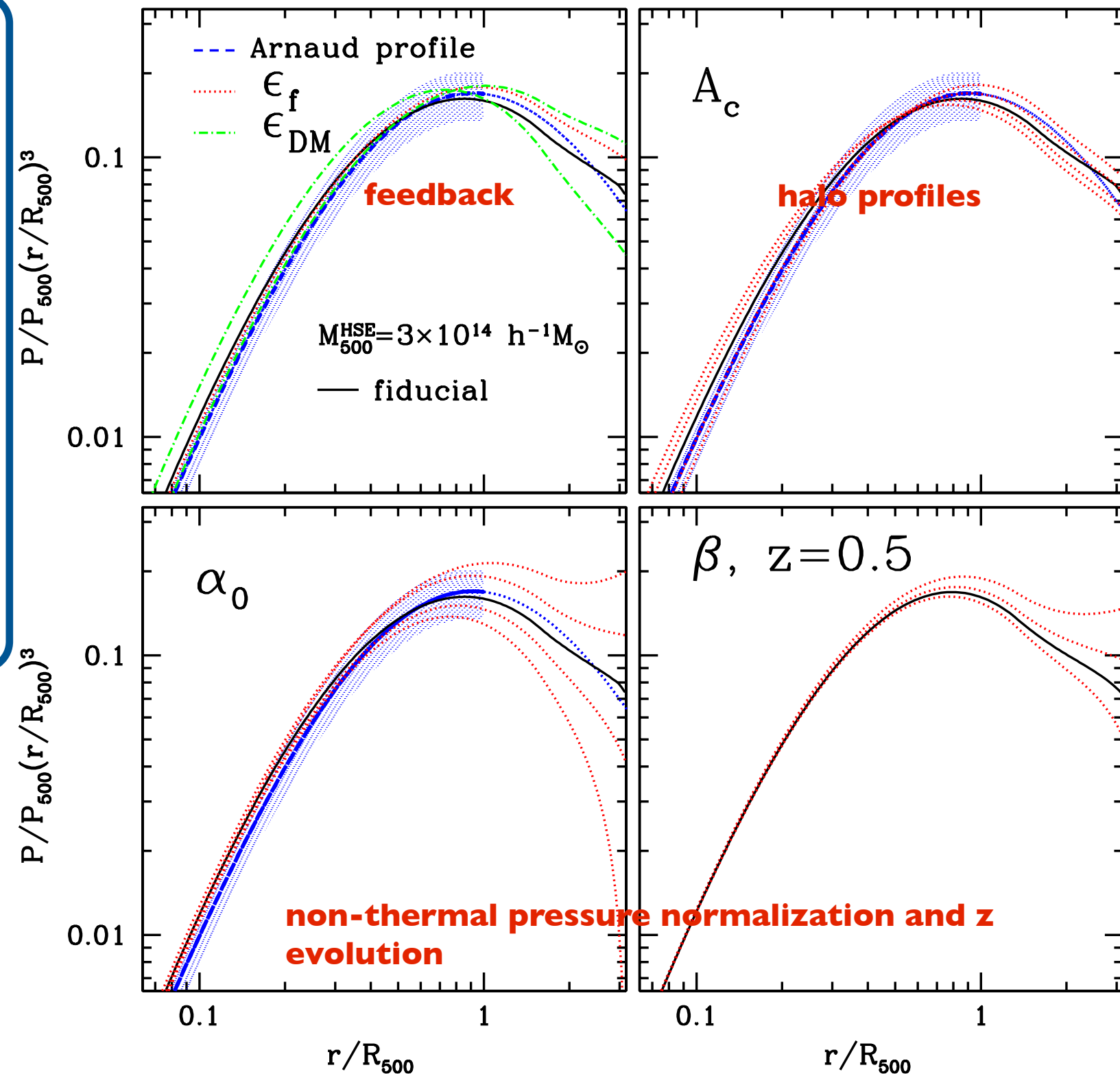
- gas reside on dark matter halos in hydrostatic equilibrium
- pressure and gas density are related via a power law
- fraction of gas turns to stars
- fraction of star energy goes back to intra-cluster medium via feedback processes.
- non-thermal pressure due to gas motion: $\frac{P_{nt}}{P_{tot}}(z) = \alpha(z) (r/R_{500})^{n_{nt}}$

Bode & Ostriker 06, ..., Shaw, Nagai, Bhattacharya & Lau 2010



Pressure Profile: Theory vs. Observation

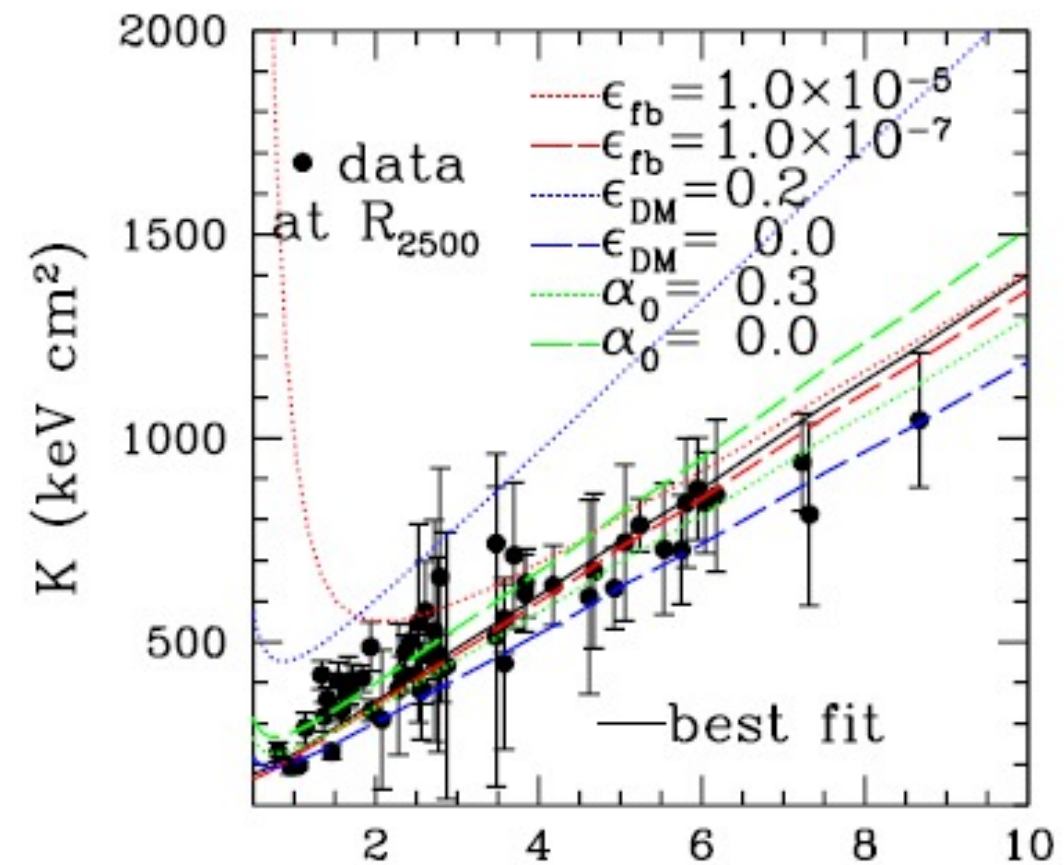
- we reproduce Arnaud profile of massive clusters at low-z -> low-z gas parameters are constrained for massive clusters
- don't have high-z pressure profile observations -> high z parameters unconstrained
- **galaxy groups -> not constrained!**



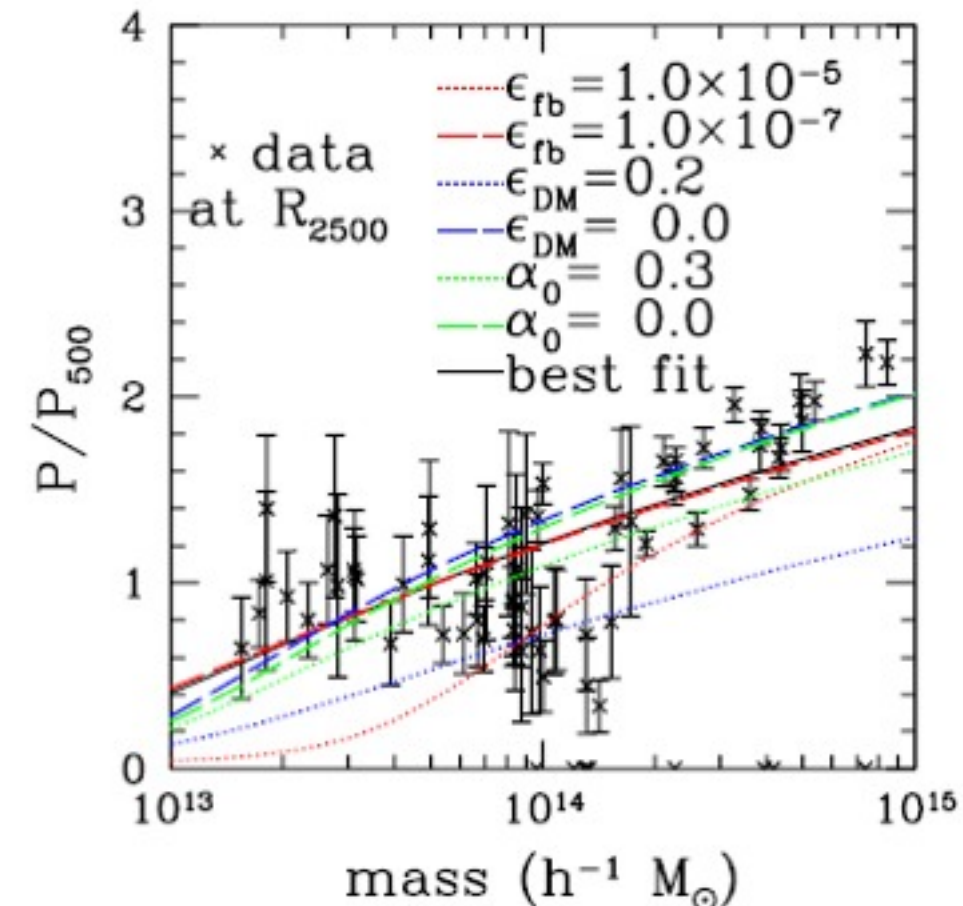
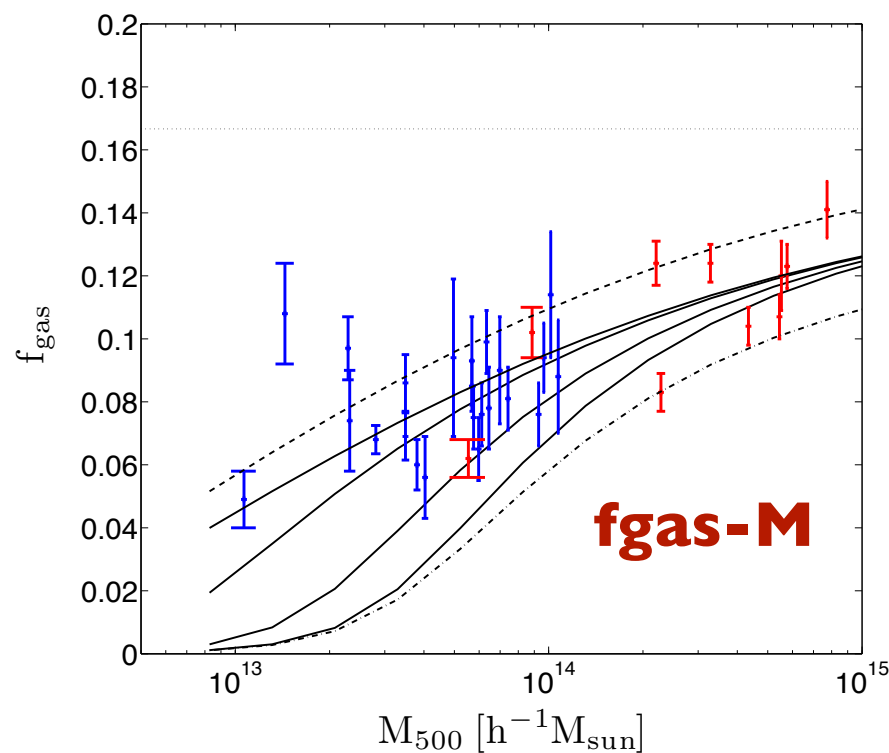
+other low-z X ray data:

K-T

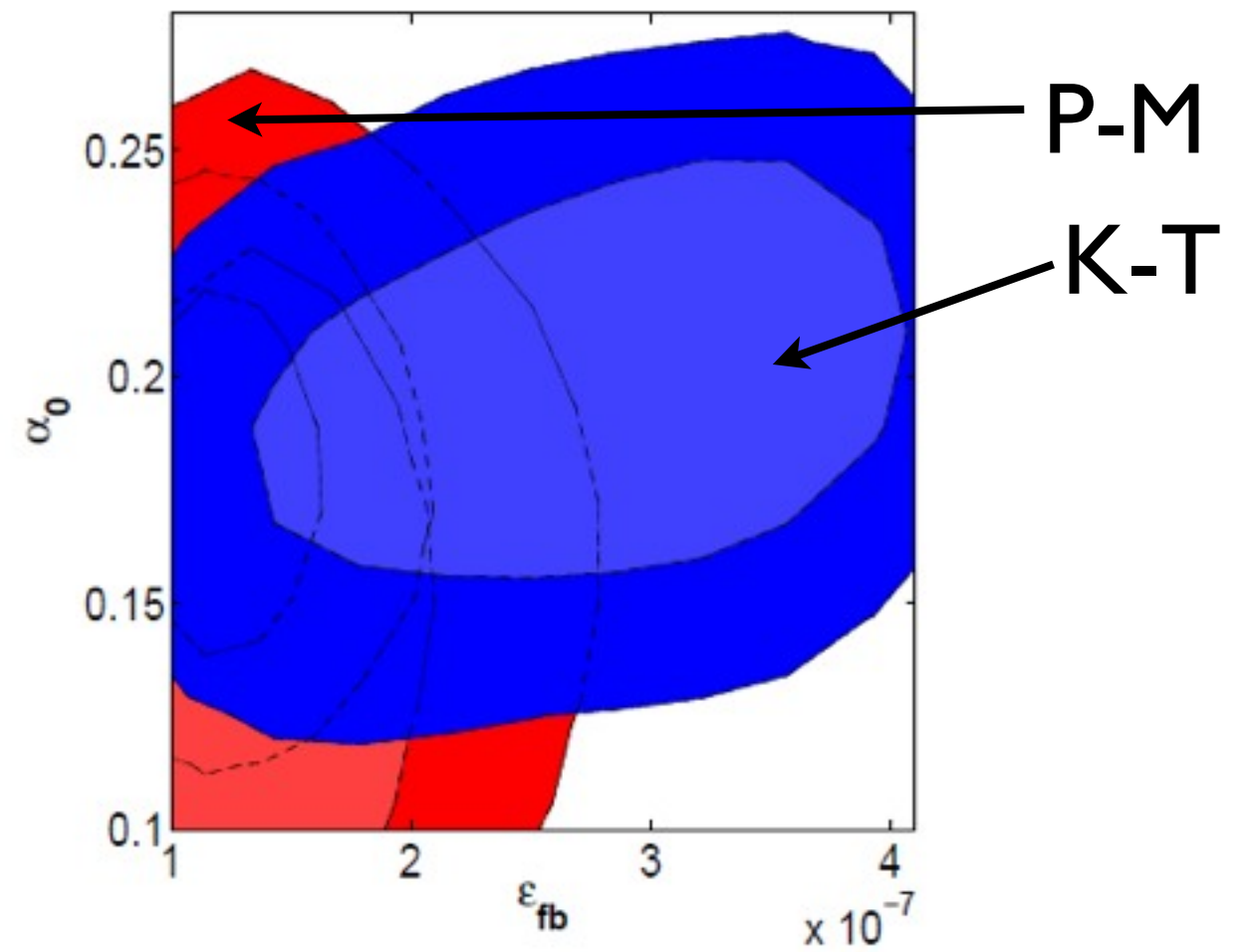
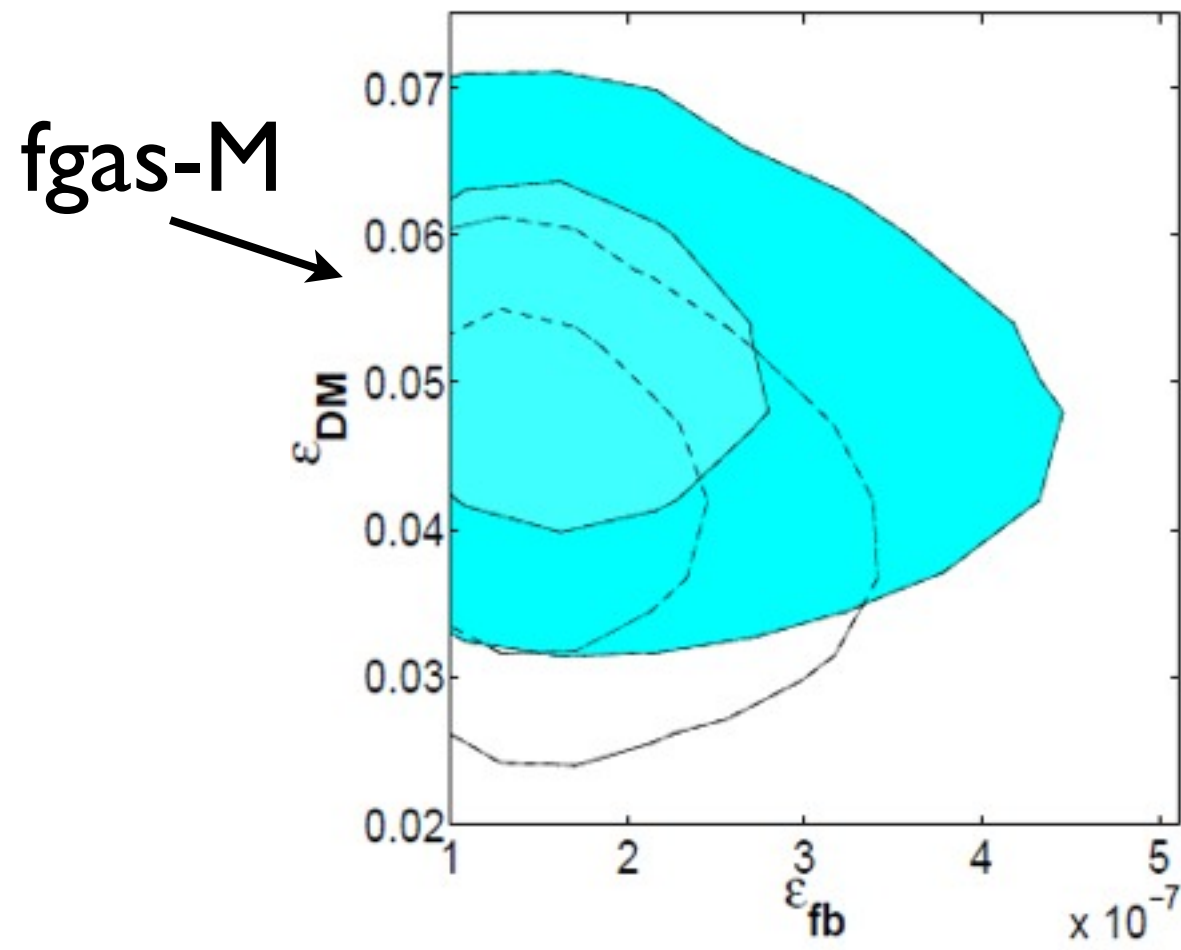
- Low redshift groups (Sun et al) and cluster (Vikhlinin et al, Pratt et al, Arnaud et al)
- **Xray scaling relations:**
 - > Vikhlinin et al , Sun et al entropy-temperature relation
 - > Pratt et al , Sun et al pressure-mass relation
 - > Vikhlinin et al , Sun et al gas fraction-mass relation



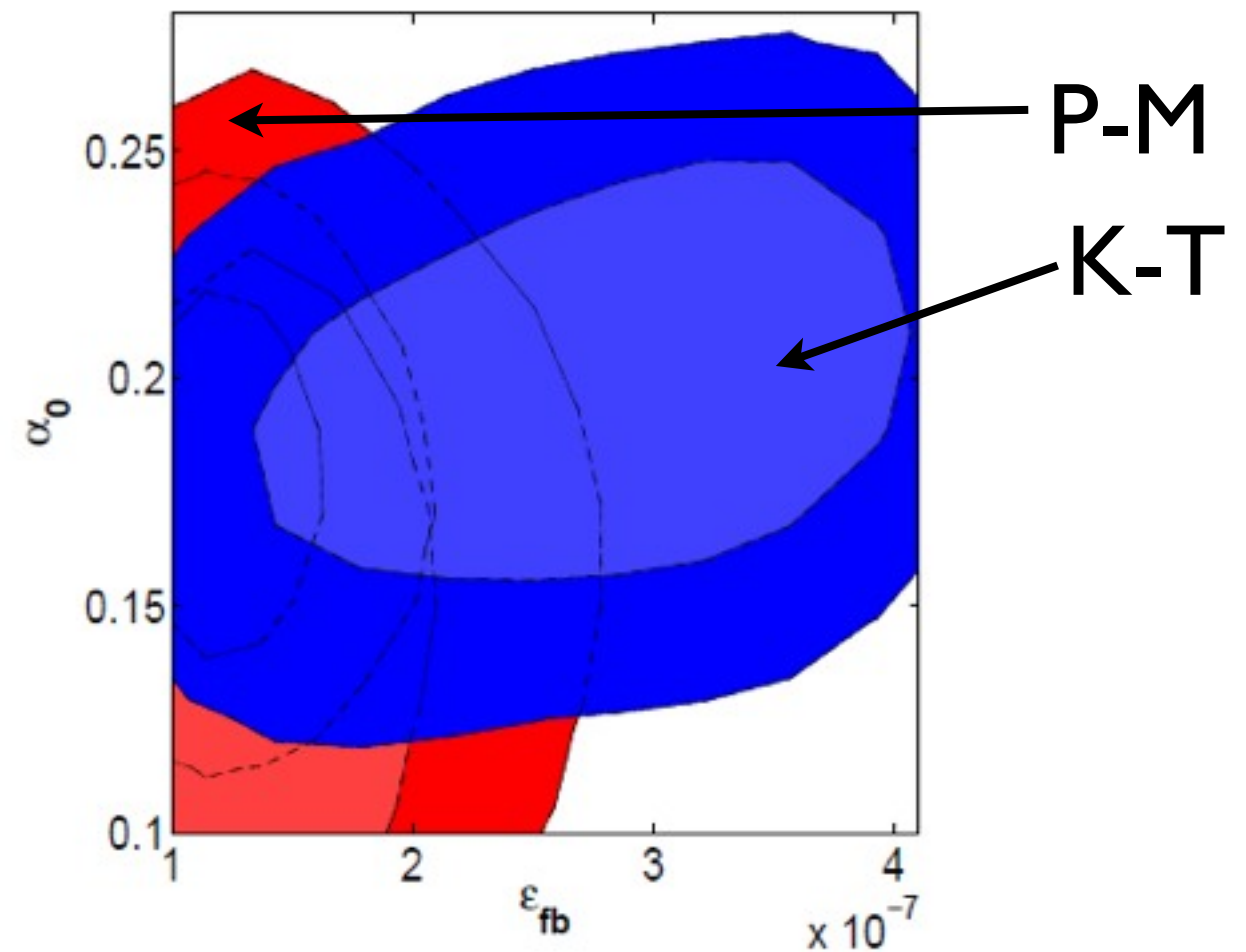
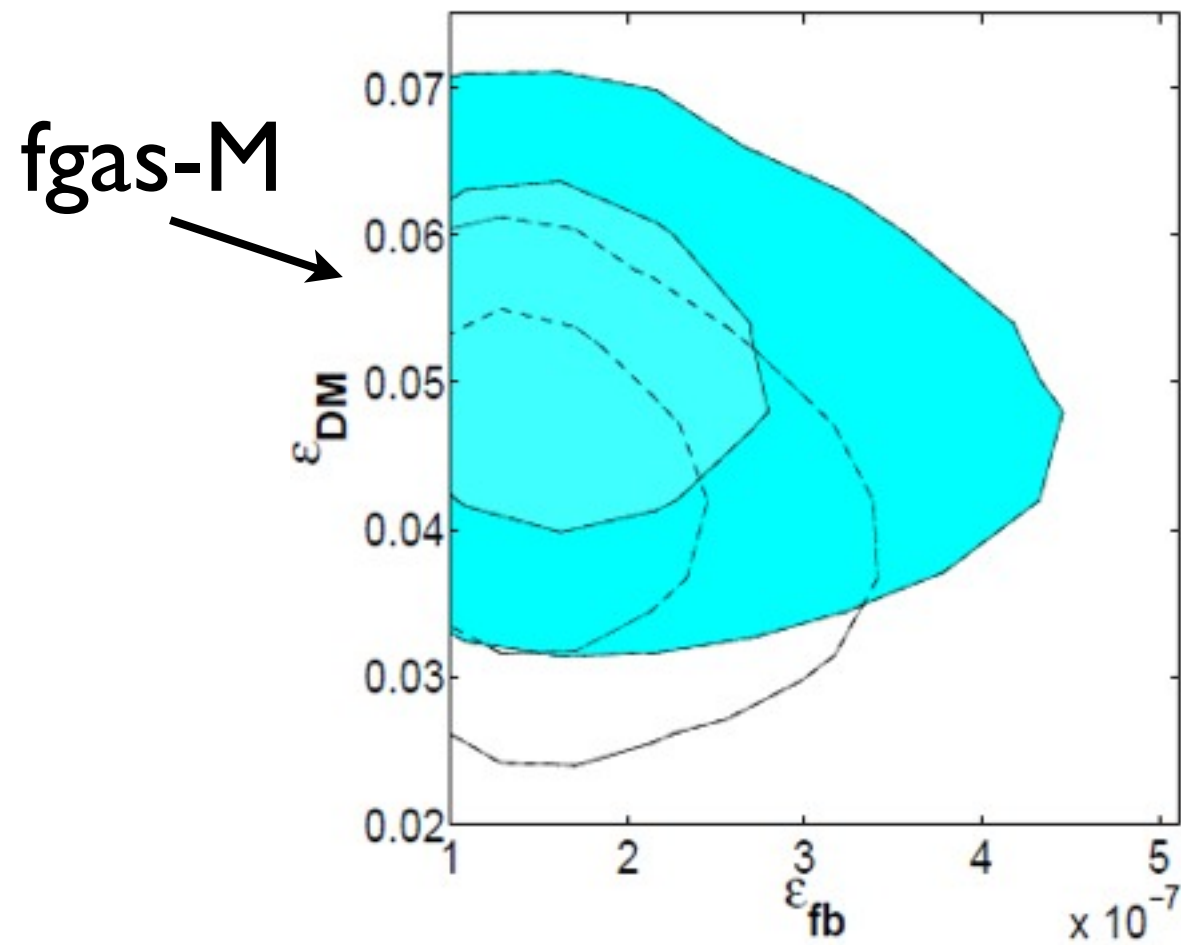
P-M



exploration of ICM parameter space



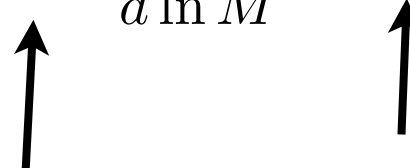
exploration of ICM parameter space



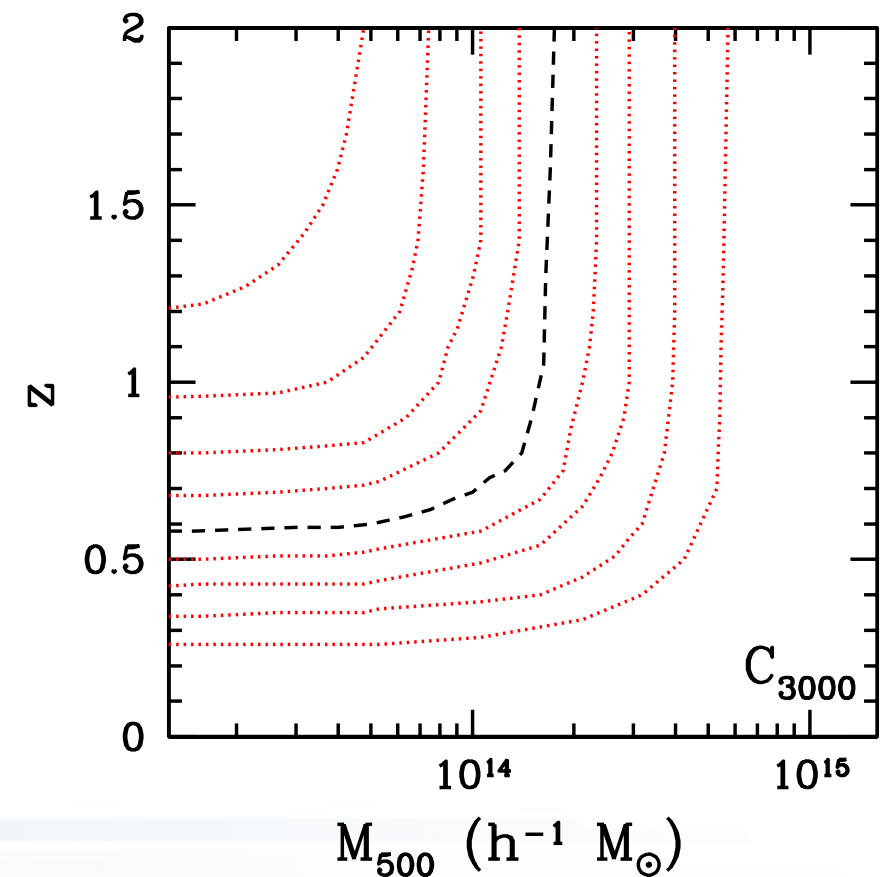
Parameters	True Mass			
	mean	std. dev.	1σ (68% CL)	2σ (95% CL)
$\epsilon_{fb} (\times 10^{-7})$	1.82	0.78	1.16-2.52	1.04-3.93
ϵ_{DM}	0.051	0.005	0.046-0.056	0.039-0.061
α_0	0.197	0.026	0.171-0.223	0.144-0.246



SZ Power Spectrum

$$C_\ell = f(x_\nu)^2 \int dz \frac{dV}{dz} \int d \ln M \frac{dn(M, z)}{d \ln M} \tilde{y}(M, z, \ell)^2$$


- $C_\ell \sim \sigma_8^8$
- **Measurement uncertainty** of the power spectrum amplitude $\sim 20\%$
- Power spectrum gets about 40% signal from high- z galaxy groups
- **Theoretical uncertainty** $\sim 40-50\%$



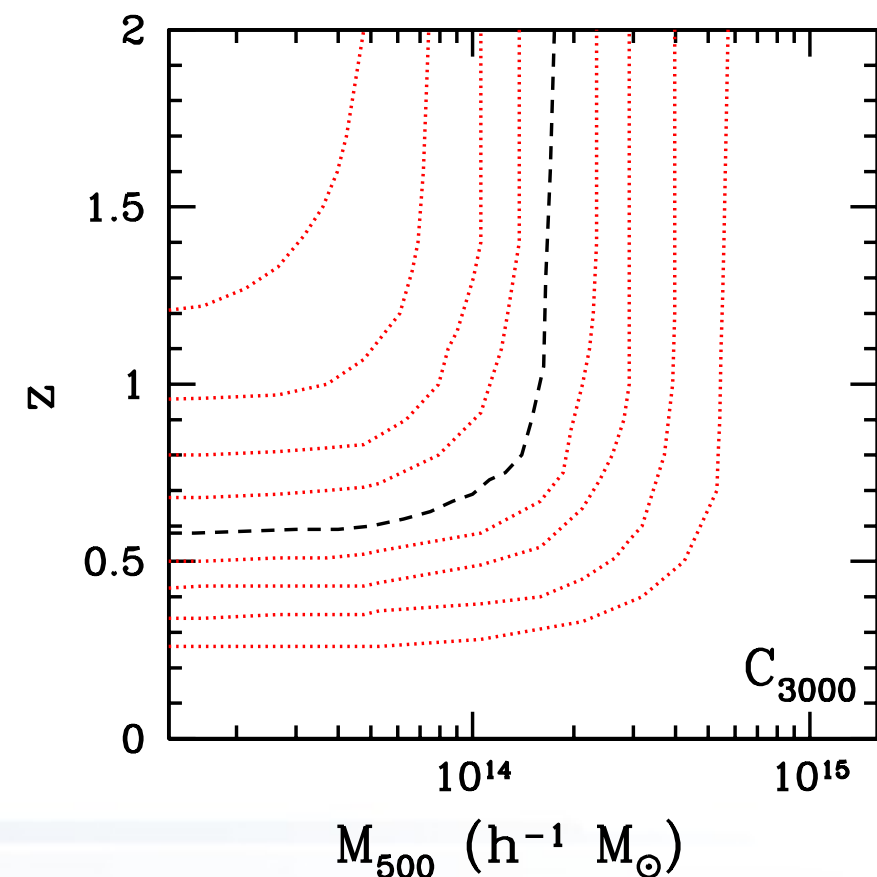
SZ Power Spectrum

$$C_\ell = f(x_\nu)^2 \int dz \frac{dV}{dz} \int d \ln M \frac{dn(M, z)}{d \ln M} \tilde{y}(M, z, \ell)^2$$

Cosmology (Geometry,
growth factor)

Astrophysics:
Fourier transform of
pressure
profile

- $C_\ell \sim \sigma_8^8$
- **Measurement uncertainty** of the power spectrum amplitude $\sim 20\%$
- Power spectrum gets about 40% signal from high-z galaxy groups
- **Theoretical uncertainty** $\sim 40-50\%$



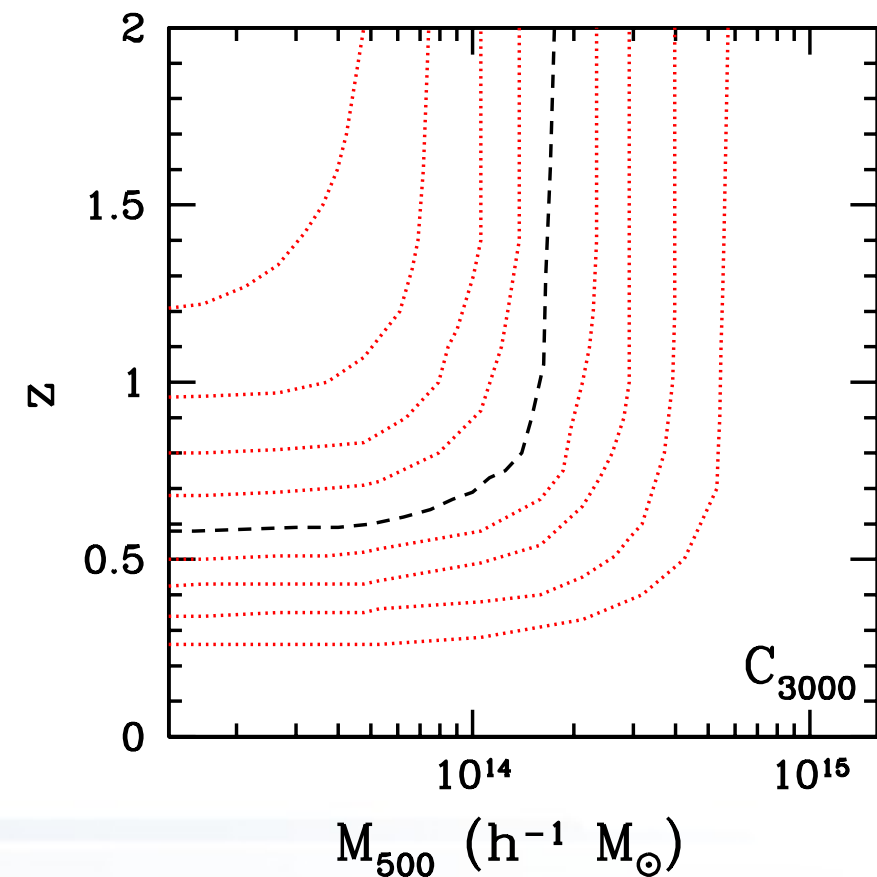
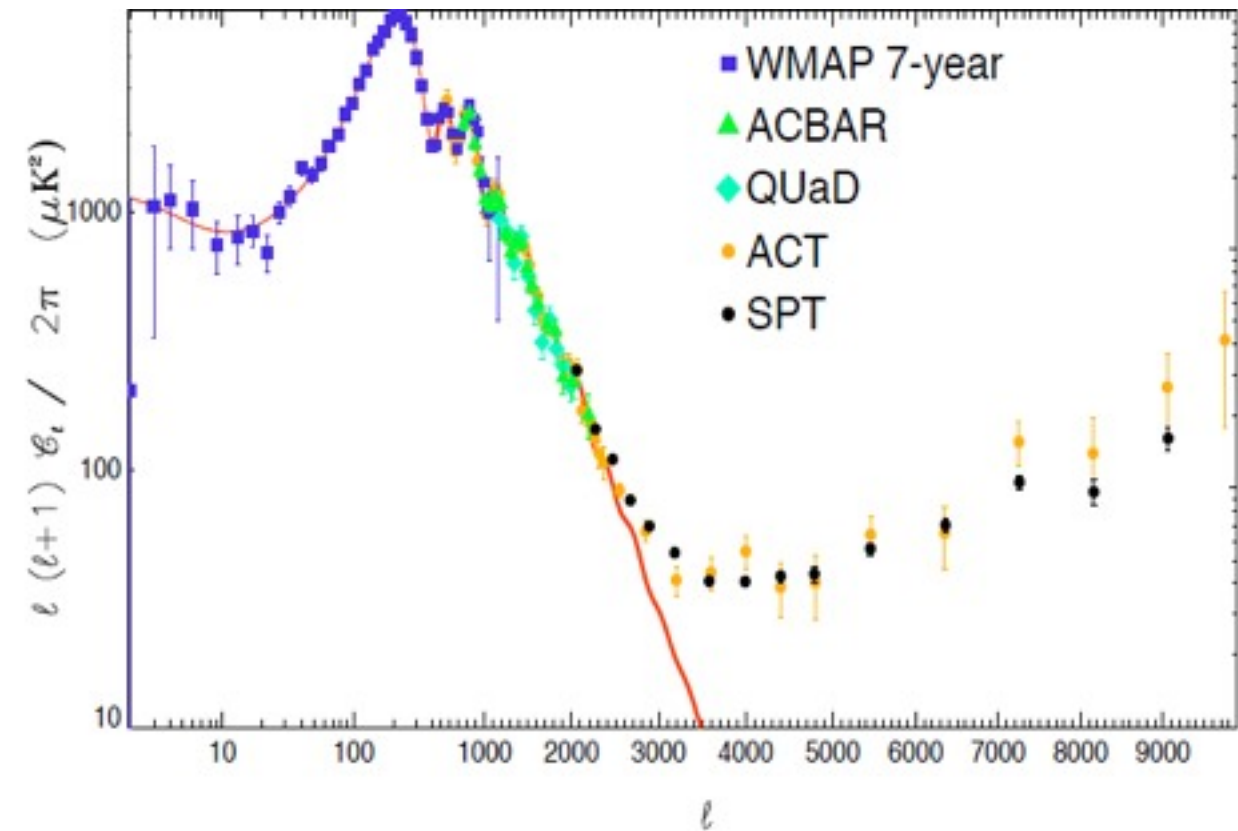
SZ Power Spectrum

$$C_\ell = f(x_\nu)^2 \int dz \frac{dV}{dz} \int d \ln M \frac{dn(M, z)}{d \ln M} \tilde{y}(M, z, \ell)^2$$

Cosmology (Geometry,
growth factor)

Astrophysics:
Fourier transform of
pressure
profile

- $C_\ell \sim \sigma_8^8$
- **Measurement uncertainty** of the power spectrum amplitude $\sim 20\%$
- Power spectrum gets about 40% signal from high-z galaxy groups
- **Theoretical uncertainty** $\sim 40-50\%$



Bispectrum Primer

- 3-pt function on the sky in Fourier space

- CMB temperature (harmonic space) in a particular direction in the sky: $a_{\ell m} = \int d^2 \hat{n} \frac{\Delta T}{T} Y_{\ell m}^*(\hat{n})$

- angular bispectrum then is: $\langle a_{\ell_1 m_1} a_{\ell_2 m_2} a_{\ell_3 m_3} \rangle$

- average over m satisfying:

$$m_1 + m_2 + m_3 = 0, \ell_1 + \ell_2 + \ell_3 = \text{even}, \text{ and } |l_i - l_j| \leq l_k \leq l_i + l_j$$

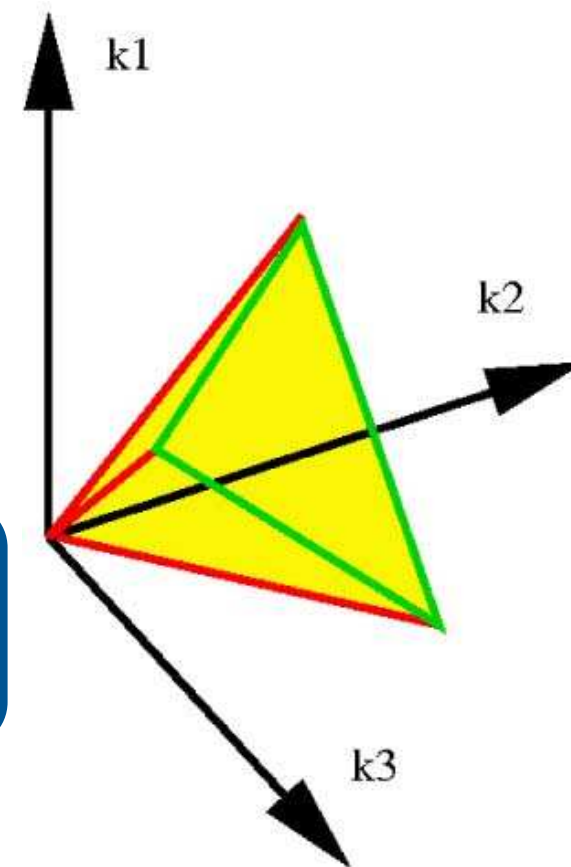
- **SZ bispectrum (flat sky limit):**

$$b(\ell_1 \ell_2 \ell_3) = f(x_\nu)^3 \int dz \frac{dV}{dz} \int d \ln M \frac{dn(M, z)}{d \ln M} \\ \times \tilde{y}(M, z, \ell_1) \tilde{y}(M, z, \ell_2) \tilde{y}(M, z, \ell_3)$$

and its noise: $N(\ell_1 \ell_2 \ell_3) = C(\ell_1) C(\ell_2) C(\ell_3)$

where $C(\ell)$ is the total power spectrum
(point sources +SZ+ lensed CMB+beam noise)

FT of the
pressure
profiles



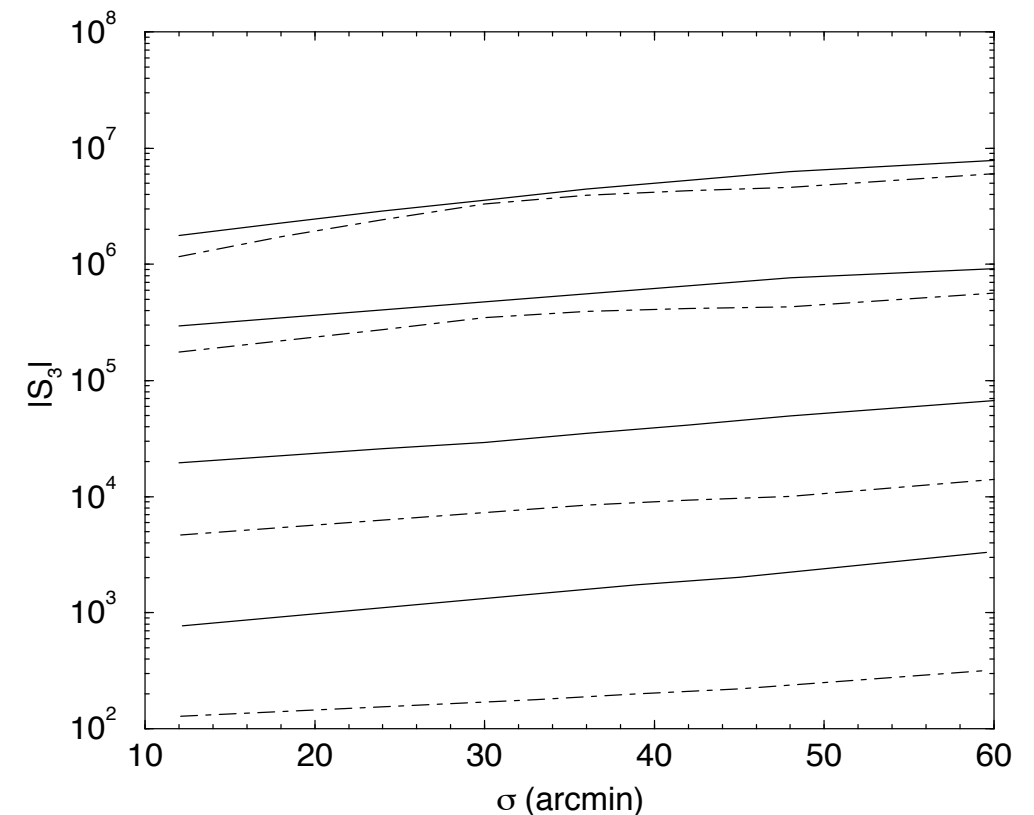
- Komatsu & Spergel 2000; Cooray & Hu 1999, 2000; Hu 2001



From Bispectrum to Skewness Spectrum

- skewness is the simplest 3-pt statistics in real space (equivalent to variance in 2-pt space):
- Skewness is the sum over all possible triangles in harmonic space, then FT to real space.
- in real space, a skewness function is the skewness measured over a certain angular scale (Cooray 2000)
- **Problem** with real space skewness function is different sources of non-Gaussianity.
- **Solution ?** define a skewness spectrum in harmonic space (Munshi & Heavens 2008)

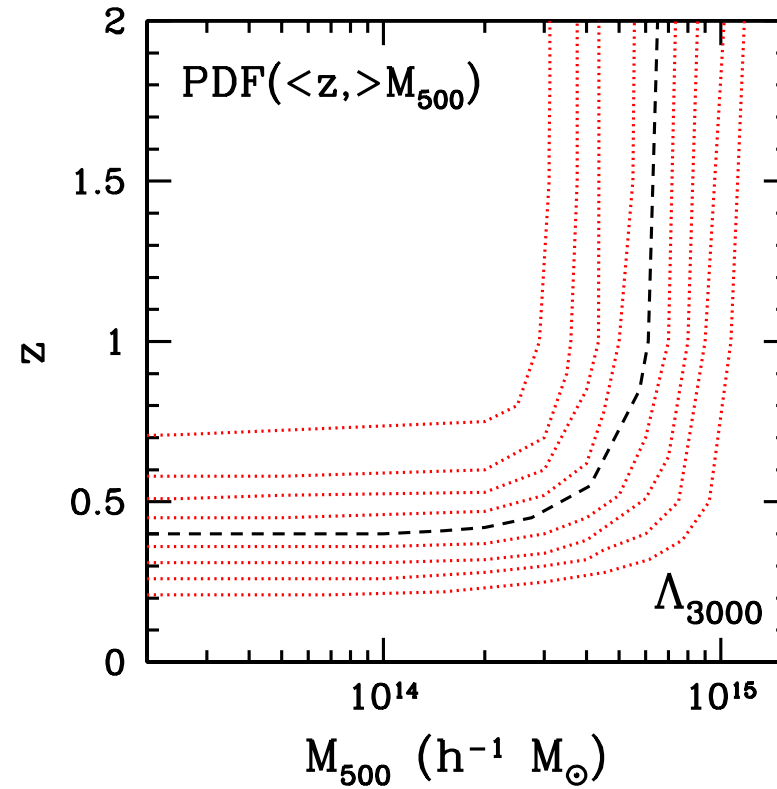
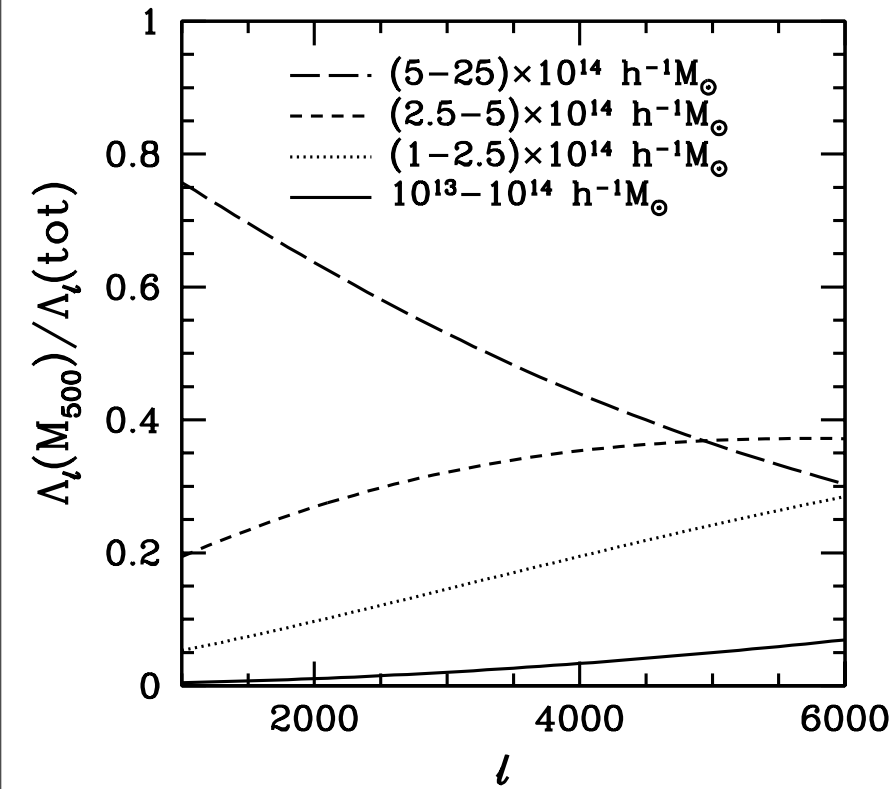
$$S_3 \equiv \left\langle \left(\frac{\Delta T(\hat{\mathbf{n}})}{T} \right)^3 \right\rangle$$



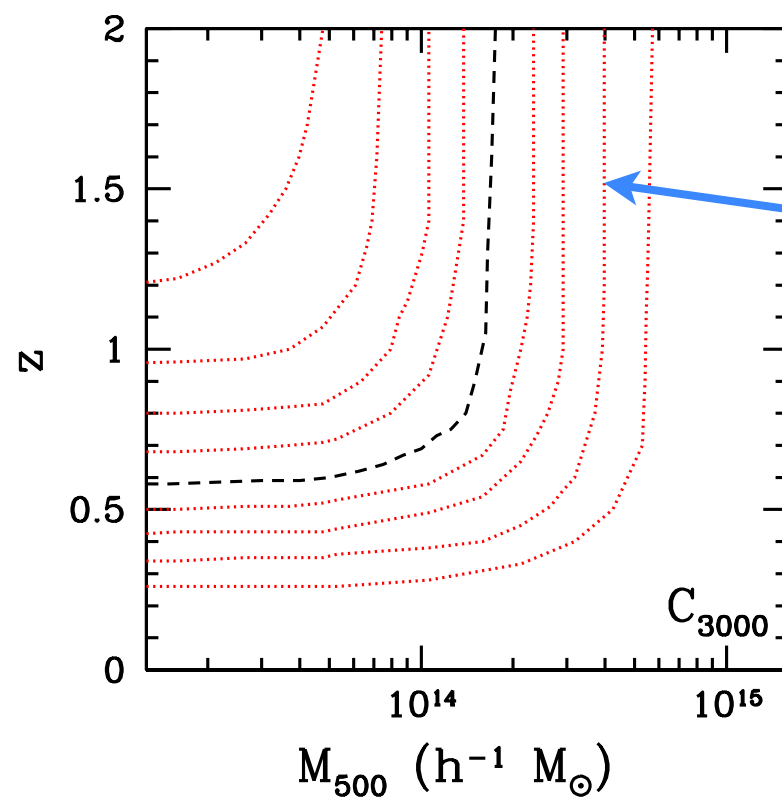
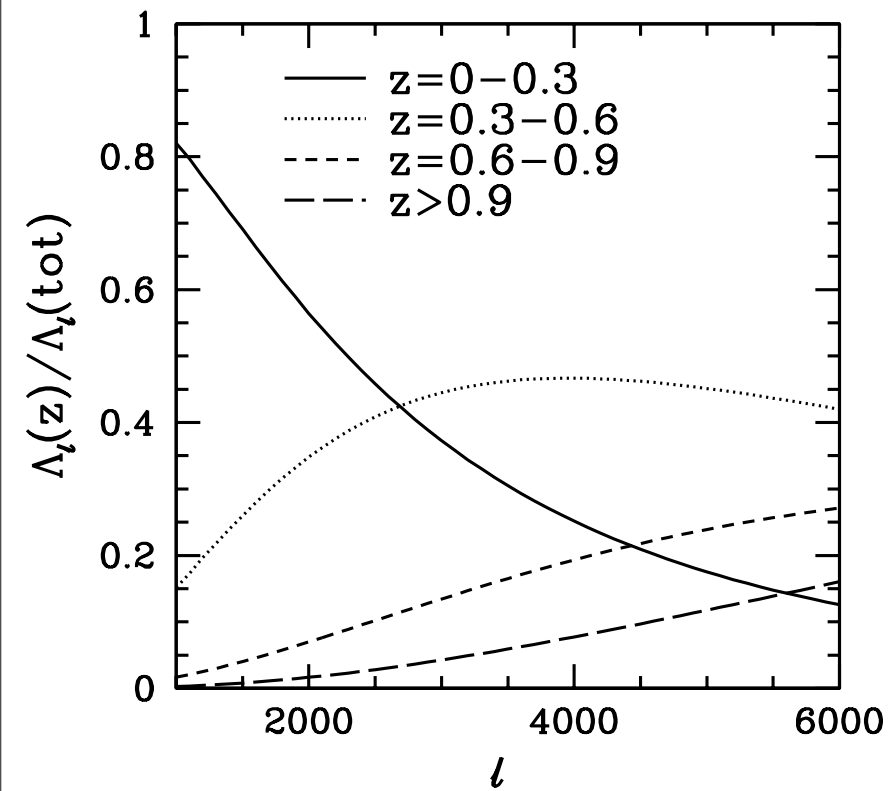
Cooray 2000; Rubino-Martin & Sunyaev 03; Hill & Sherwin 12



Where the SZ skewness spectrum signal comes from?



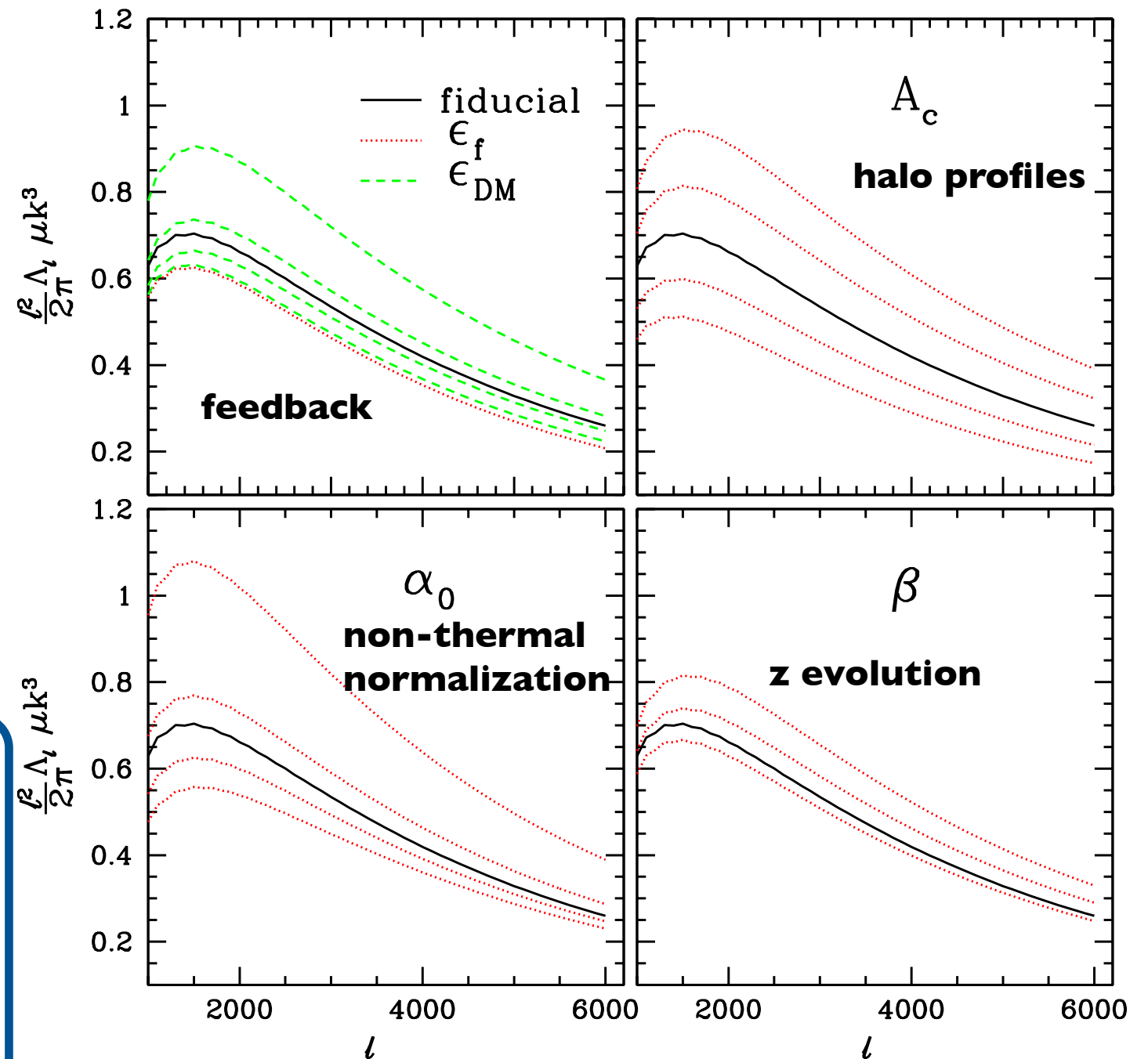
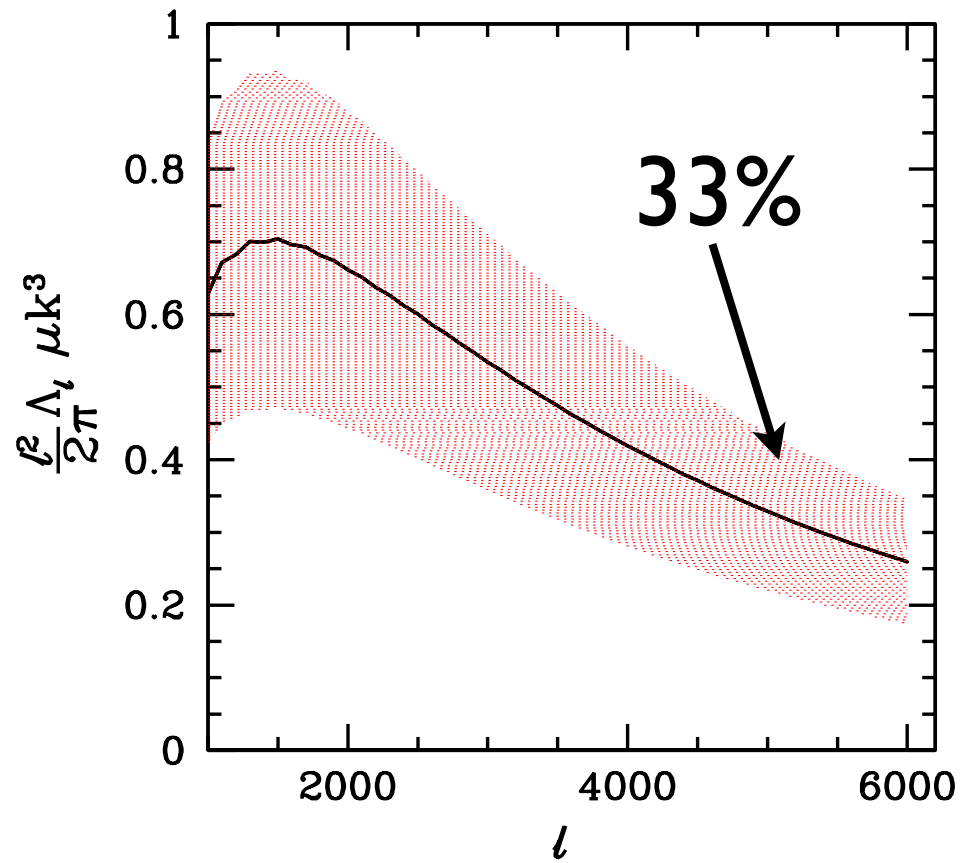
about 95% of the signal comes from $M_{500} > 1e14$ & $z < 1$



SZ power spectrum distribution (also Trac et al 2011)



Astrophysical Uncertainties of the Skewness Spectrum:



- theoretical uncertainty in the skewness spectrum $\sim 33\%$ (compared to the 40-50% uncertainty in the SZ power spectrum)



Power Spectrum vs. Bispectrum

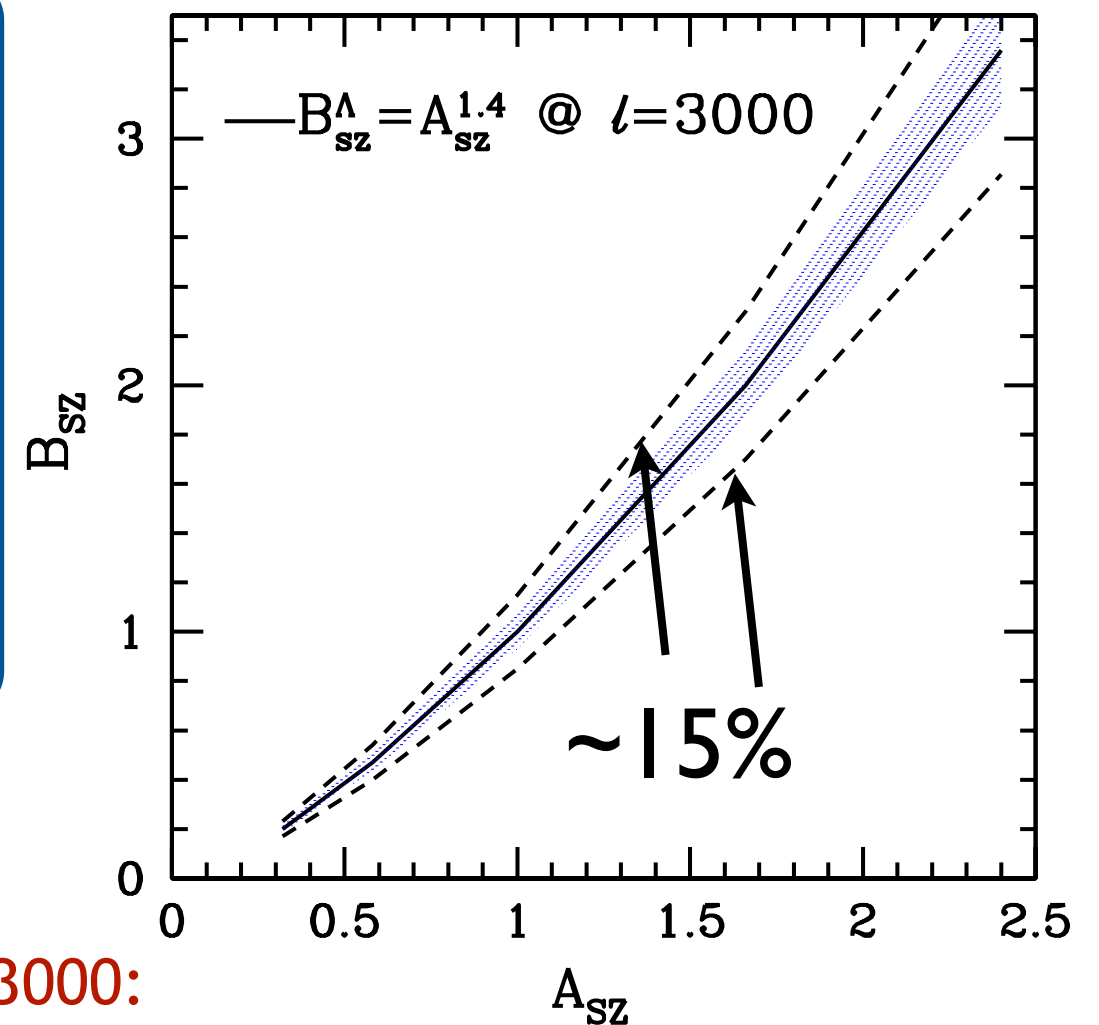
- SZ power spectrum gets about 30-50% signal from groups and high-z clusters vs. bispectrum gets only 5% signal from such objects.
- power spectrum measures a combination of kinetic+thermal SZ (and point sources) vs. kinetic SZ bispectrum =0, so bispectrum measures thermal SZ only.

- define the SZ skewness spectrum amplitude @ $l=3000$: $B_{sz} = \Lambda_{4000}(\sigma_8) / \Lambda_{4000}(0.8)$

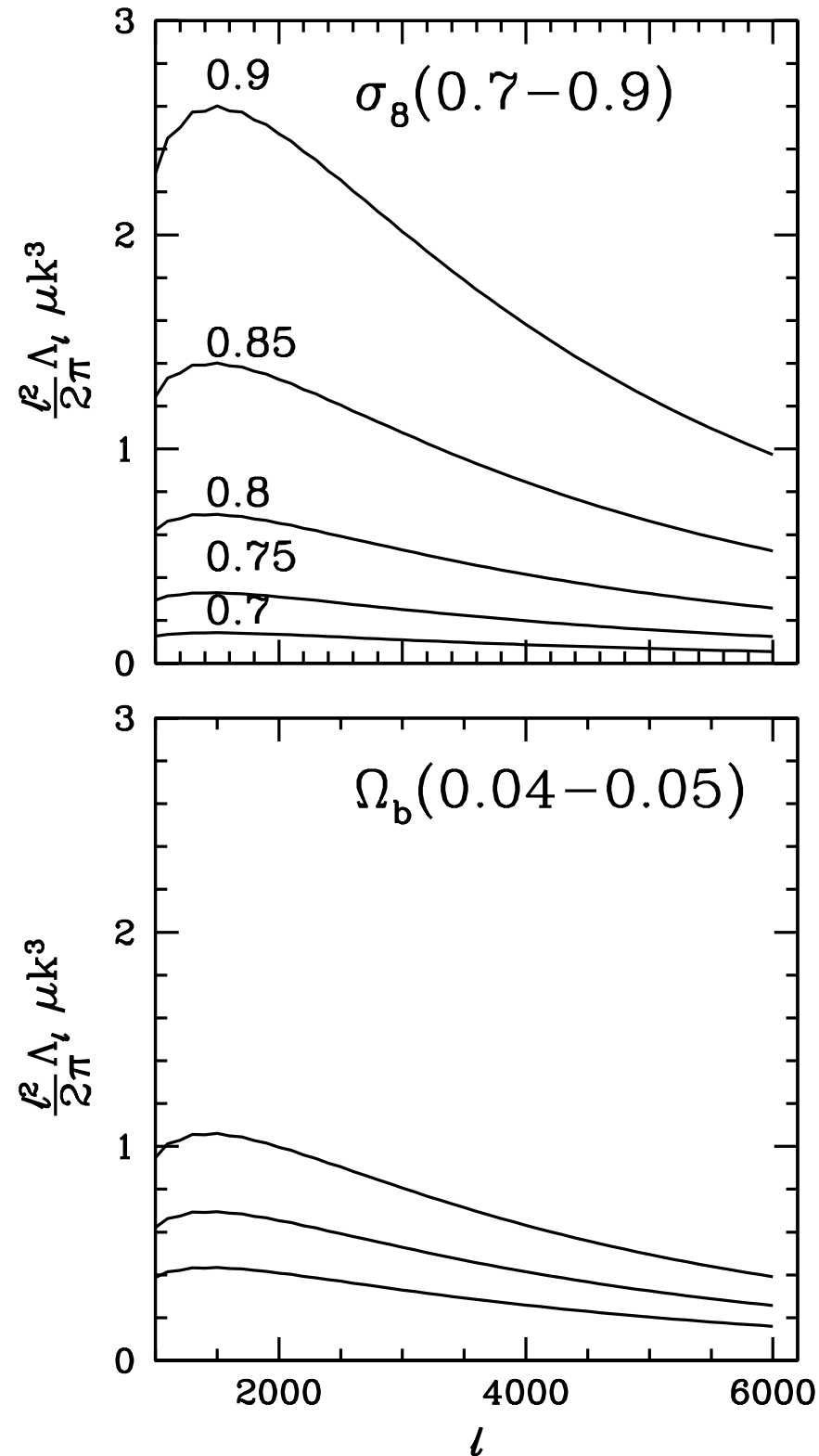
- define the SZ power spectrum amplitude @ $l=3000$:

$$A_{sz} = C_{3000}(\sigma_8) / C_{3000}(0.8)$$

- Asz-Bsz relation is extremely robust. Change in gas physics changes it only by 15%. The power spectrum over the same range changes by factor of 4.
- A combination of Asz and Bsz amplitude can break the degeneracy of the thermal and the kinetic SZ amplitude.



Cosmology dependence of SZ Bispectrum



Scaling of the bispectrum amplitude

$$B_{SZ}^{\Delta} \propto \left(\frac{\sigma_8}{0.8}\right)^{11.4} \left(\frac{\Omega_b}{0.04}\right)^4 \left(\frac{h}{0.71}\right)^2 \left(\frac{w_0}{-1.0}\right)^{-0.95} \\ \times \left(\frac{n_s}{0.96}\right)^{-1.5} \left(\frac{\Omega_m}{0.26}\right)^{-0.46}$$

Cosmic Complementarity->

=> SZ bispectrum constrains $\sigma_8 \Omega_b^{0.36}$
 complementary to clusters which
 constrains $\sigma_8 \Omega_m^{0.4}$

=> mask out the clusters used in the
 SZ mass function:
 => measure bispectrum of the map=>
 constrain σ_8
 => independent of the cluster
 constraints
 => joint constraints from bispectrum
 +abundance



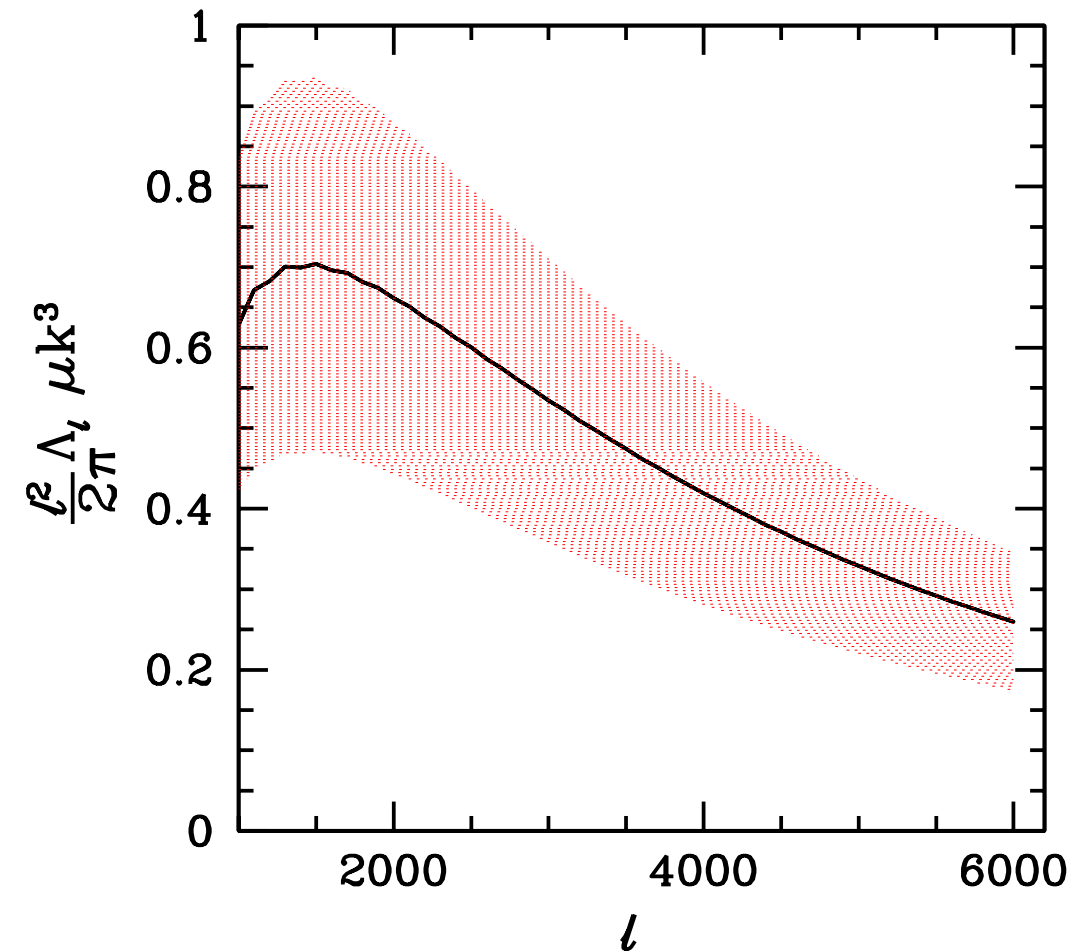
The SZ Skewness Spectrum

- define the SZ skewness spectrum as sum over the two smaller sides and expressed as a function of the largest l :

$$\Lambda(l) = \sqrt{\sum_{l_1 l_2} b^2(l l_1 l_2)}.$$

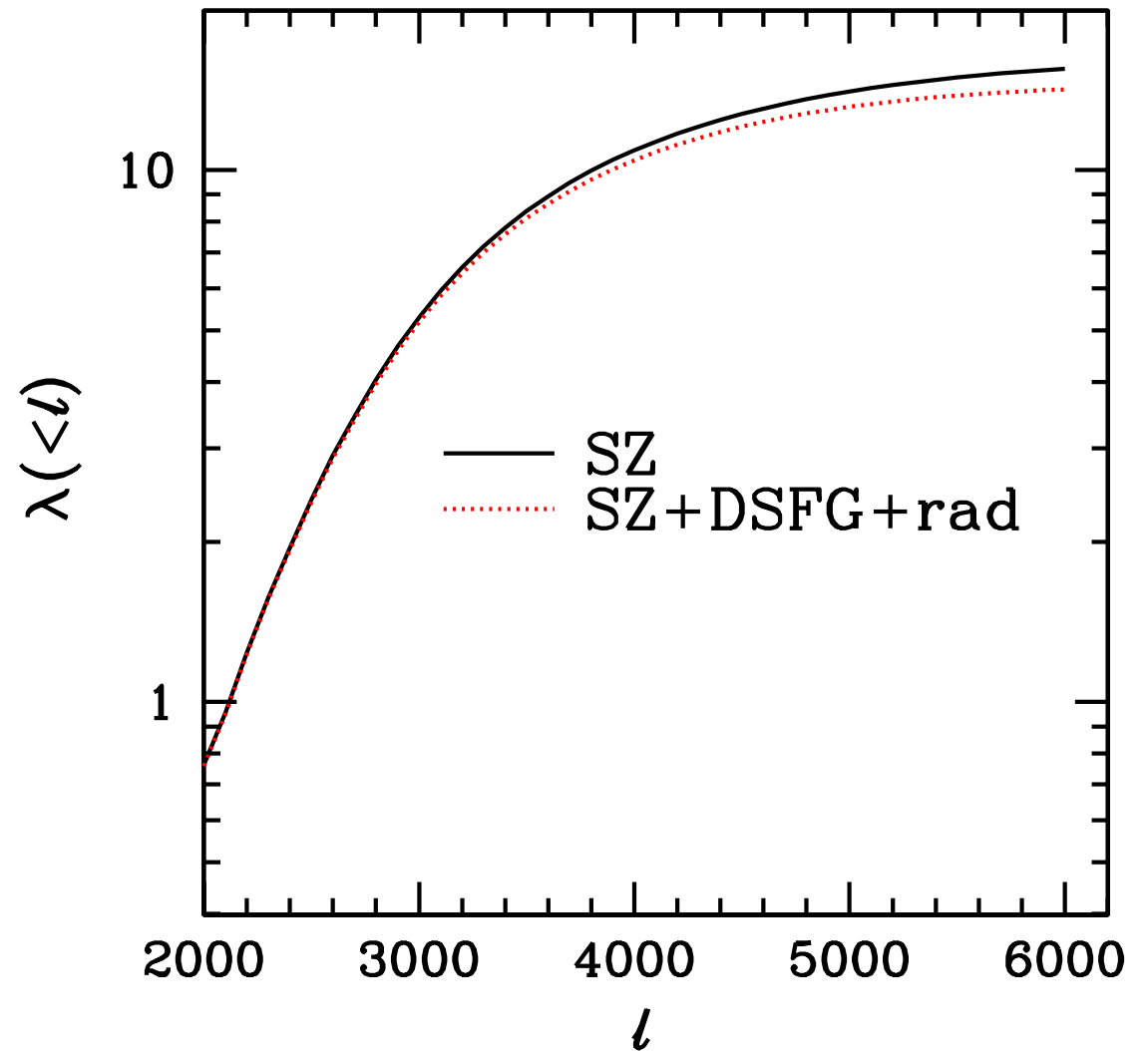
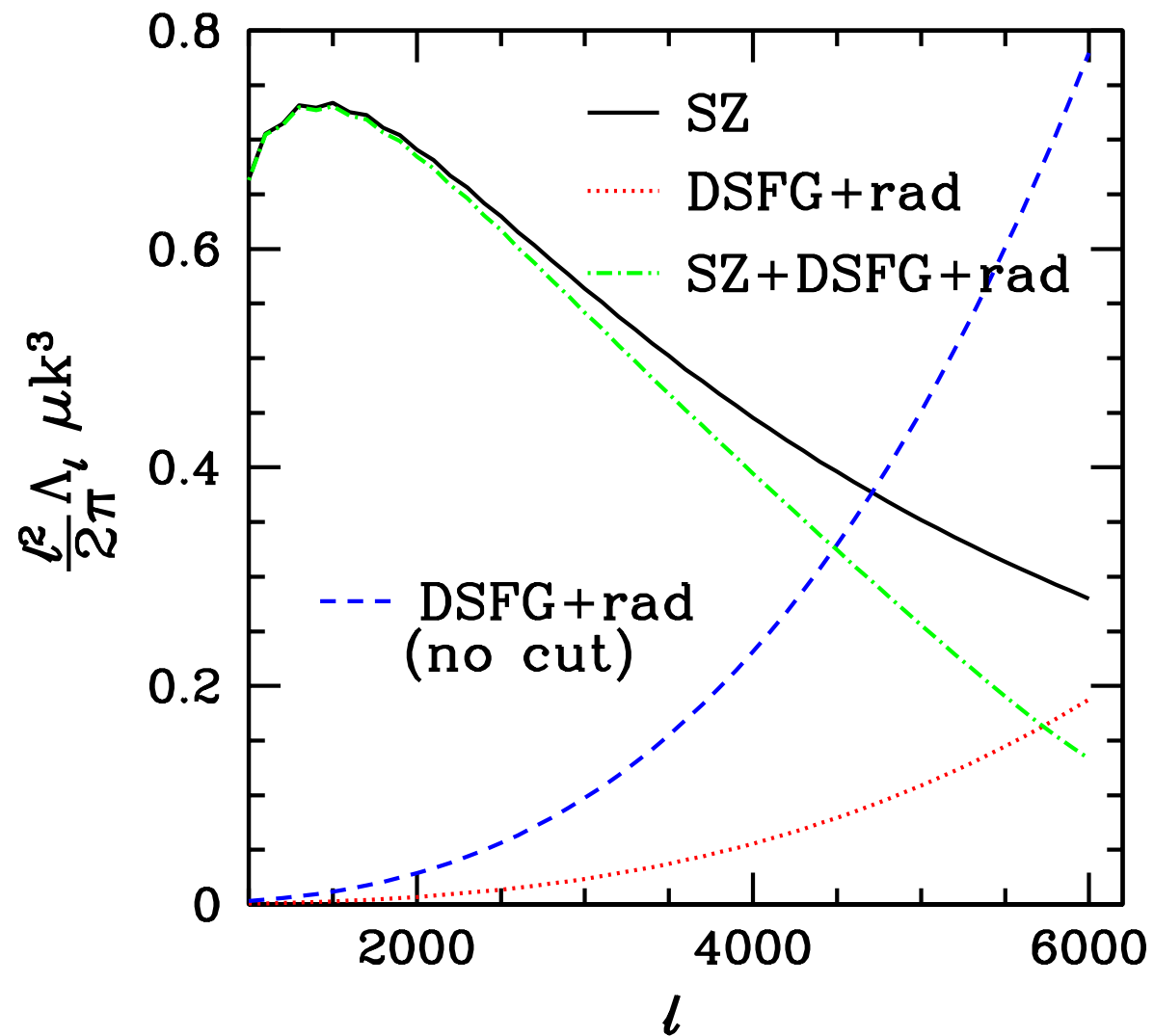
- and signal-to-noise integrated to certain l :

$$\lambda(< l) = \sqrt{\sum_{l_1}^l \sum_{l_2 l_3} \frac{b^2(l_1 l_2 l_3)}{N^2(l_1 l_2 l_3)}}$$

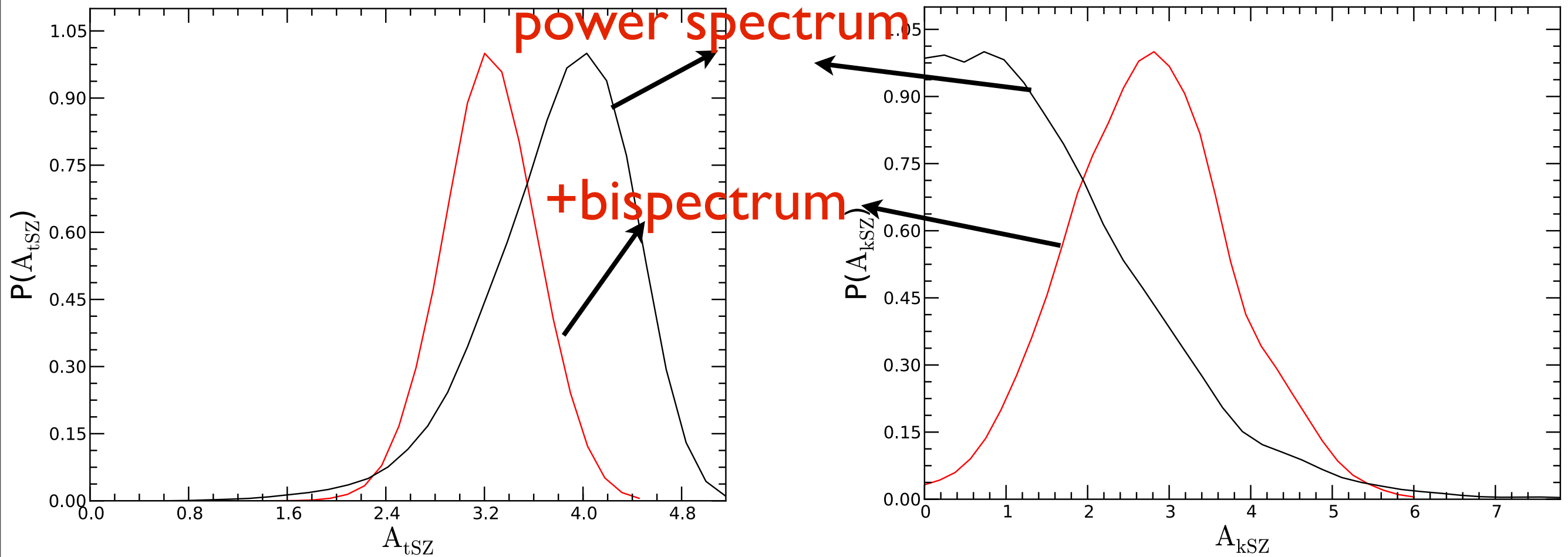


Measurement Prospects From Current Data

- @150 GHz, 1.2', 18 microk-arcmin, 2500 deg²
- total S/N~16
- Bispectrum of point sources (dusty and radio galaxies) are a contamination (or signal to detect!)



What can we learn from bispectrum+ power spectrum combined



- Adding bispectrum to the power spectrum data improves the constraints on the thermal SZ amplitude by factor of 2.
- kinetic SZ amplitude can be detected at 2 sigma level

ICM parameter constraints from future experiments

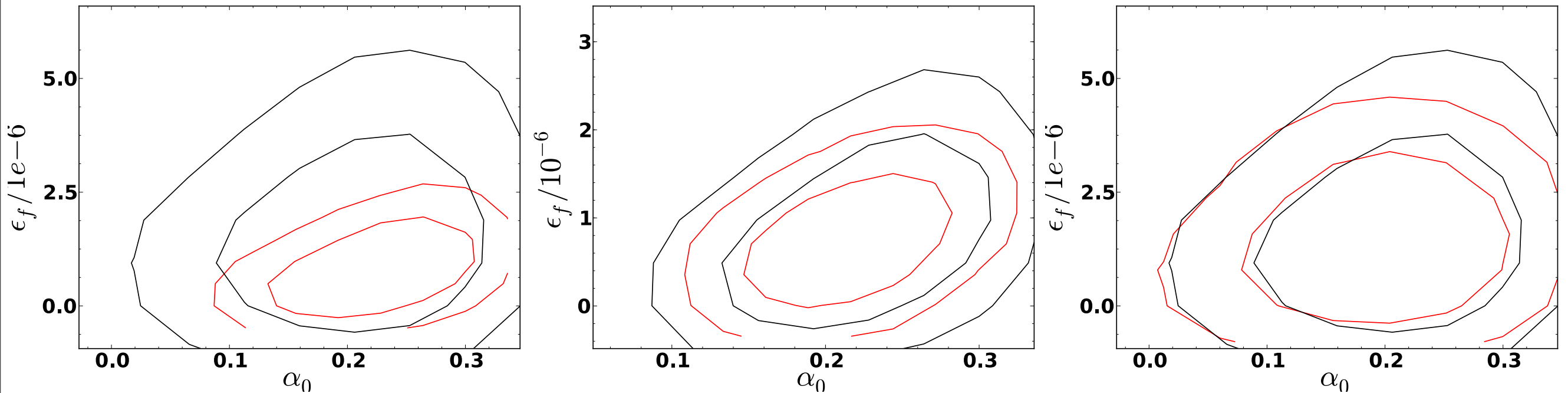
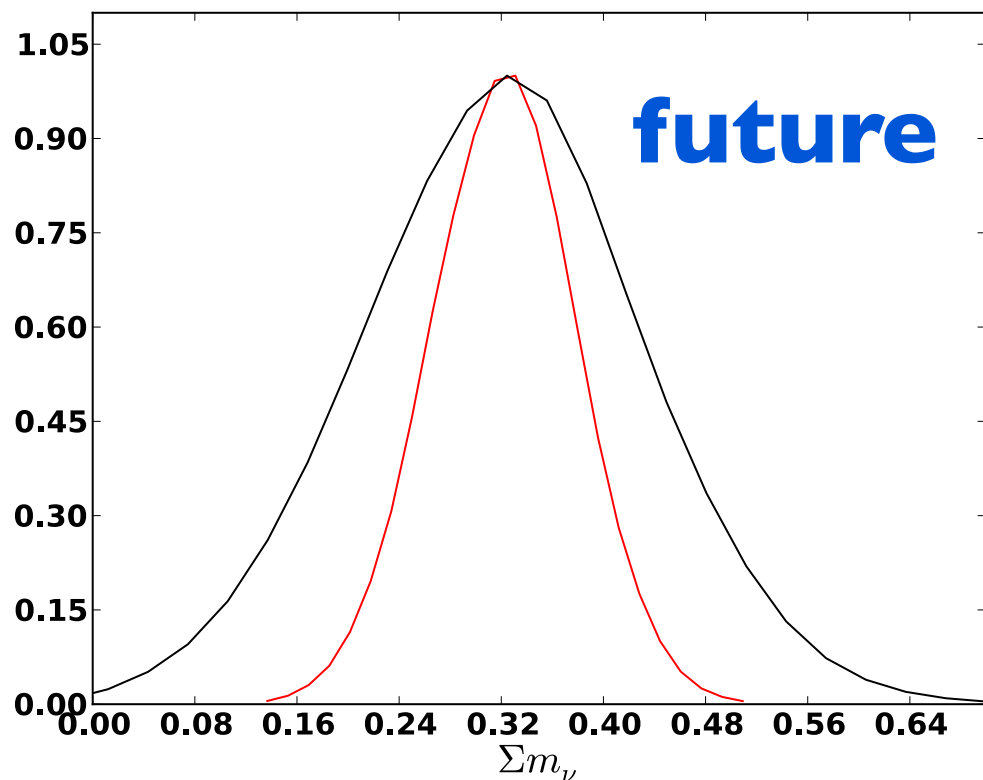
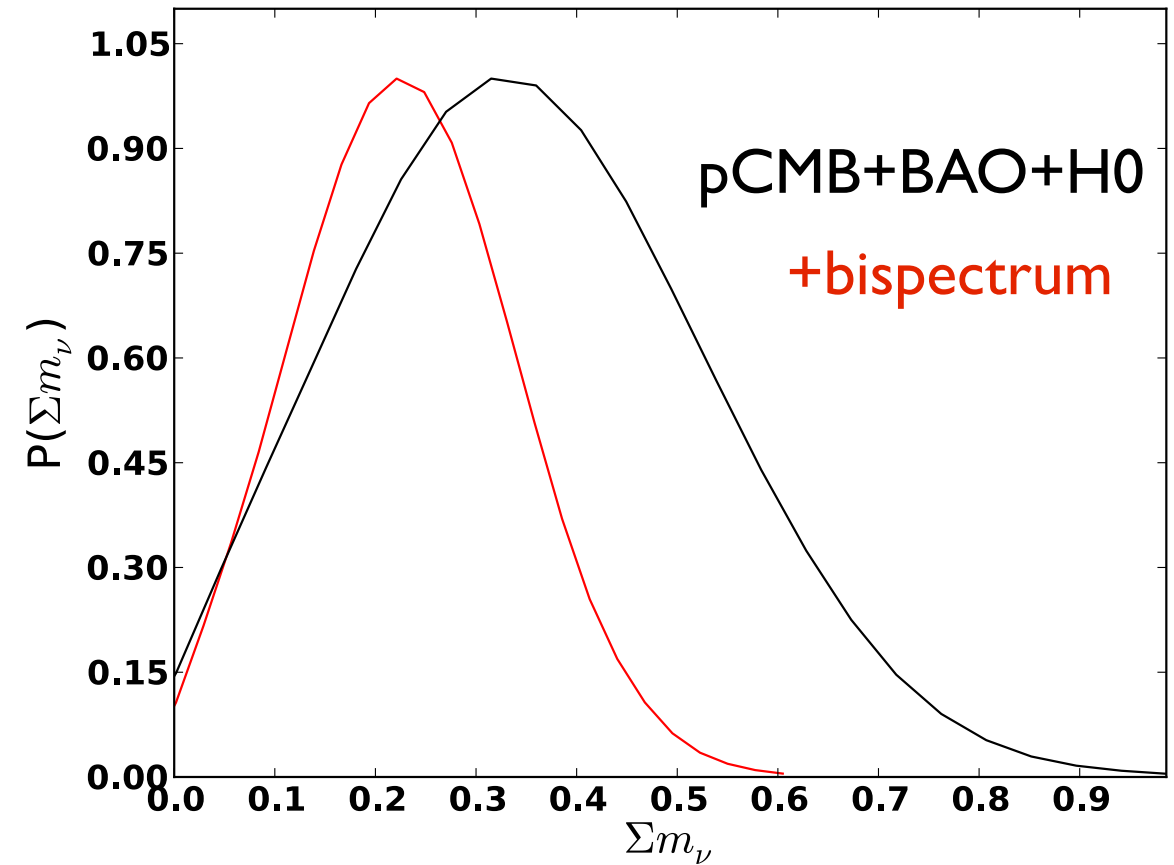
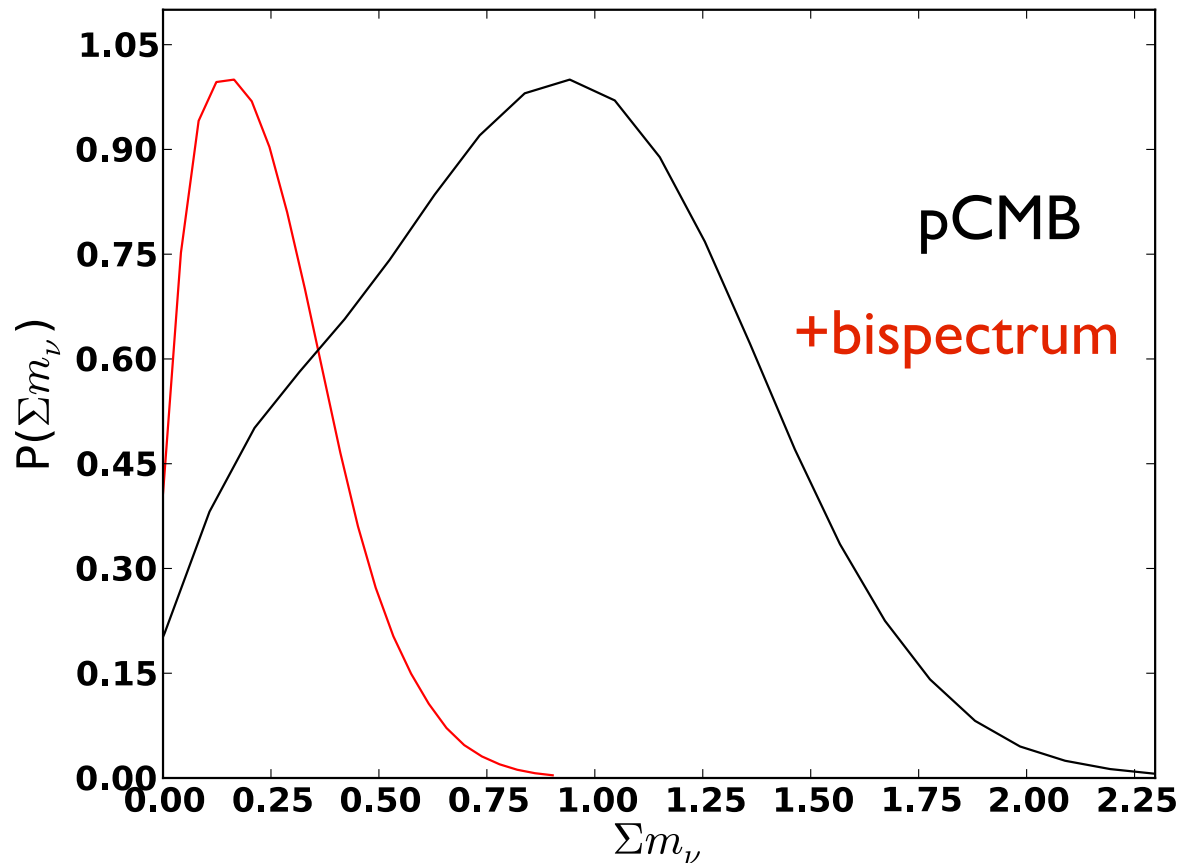


Table 5: Parameter constraints (1 (2) σ limits) from MCMC analysis of the power spectrum measurements

	CCAT (1 Khr, 2000 deg ²)		SUPERCCAT (100 Khr, 20,000 deg ²)	
survey	α_0	$\epsilon_f(10^{-6})$	α_0	$\epsilon_f(10^{-6})$
fix cosmo, medsz, fg	0.15– 0.23(0.13 – 0.29)	1.35– 4.45(0.36 – 5.8)	0.18–0.25(0.15 – 0.29)	0.34–1.17(0 – 1.58)
w7, medsz, fg	0.07– 0.25 (0.015– 0.29)	0.86– 3.8 (0.2 – 5.65)	0.17– 0.27(0.12 – 0.29)	0.24– 1.41(0 – 2.12)
w7, fg	0.073–0.25 (0–0.29)	1.6– 6.2(0.32 – 8.6)	0.14–0.27(0.07 –0.29)	0.4–2.9(0 – 4.56)
w7, nofg	0.075– 0.25(0.02 – 0.29)	0.67–4.0 (0.11 – 6.2)	0.13– 0.25 (0.067 – 0.29)	0.45– 2.5(0.07–3.61)



Constraints on the sum of the neutrino mass



adding bispectrum

-> improves neutrino mass constraints by factor of 3 compared to WMAP alone

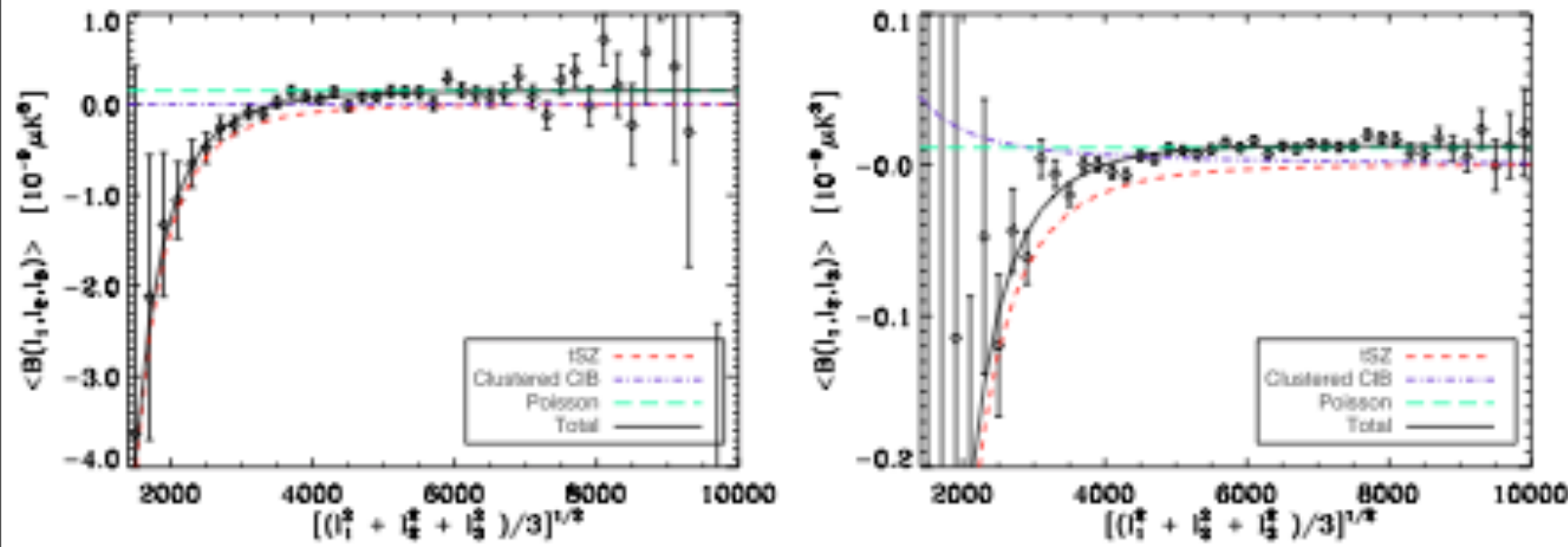
-> and by about 50% compared to WMAP +BAO+H0

-> Future: adding bispectrum can constrain neutrino mass with 0.06-0.1 eV accuracy.



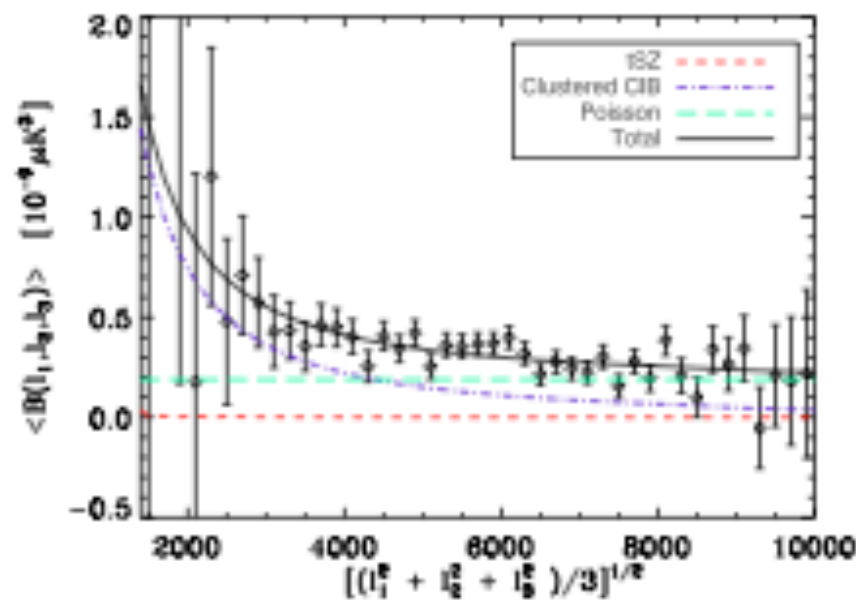
SPT Bispectrum Measurements

Crawford et al, SPT team

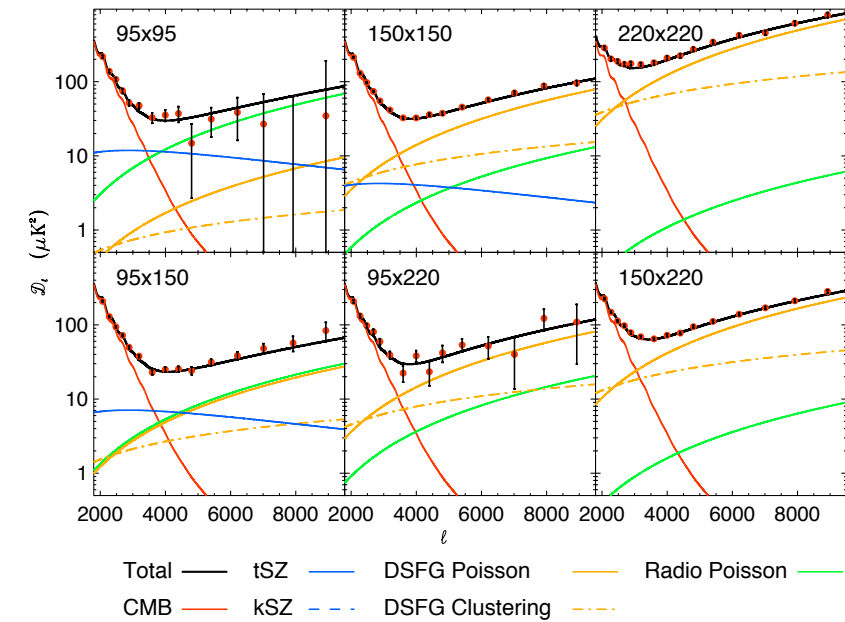


(a) 95 GHz 1d bispectrum, with best-fit model overplotted

(b) 150 GHz 1d bispectrum, with best-fit model overplotted



(c) 220 GHz 1d bispectrum, with best-fit model overplotted



- => SPT 3 frequency channels- 95, 150, 220 GHz
- => cover 800 sq. deg
- => detect SZ bispectrum to $> 10 \sigma$

Also, Wilson et al, ACT measure skewness $\sim 5 \sigma$
 Planck measurements of SZ bispectrum



Current SPT measurements: thermal SZ and σ_8

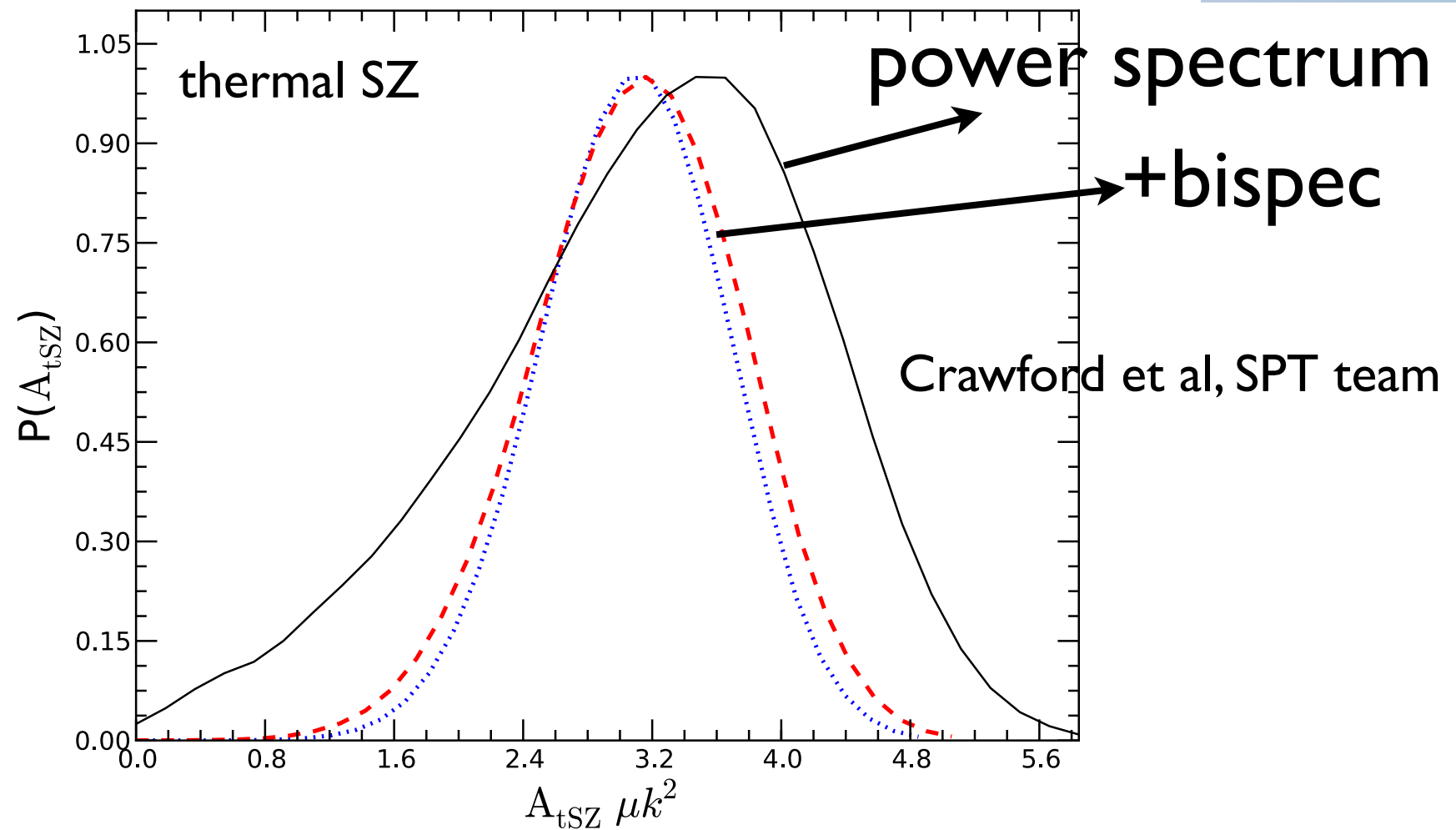
two theoretical modeling assumptions

10% (nominal) and 18% (extreme)

nominal- one that fits the X-ray data

extreme- one that test the physical boundary of our model

power spectrum changes by factor ~ 4 in the extreme scenario



=> bispectrum measures $\sigma_8 = 0.79 \pm 0.03$

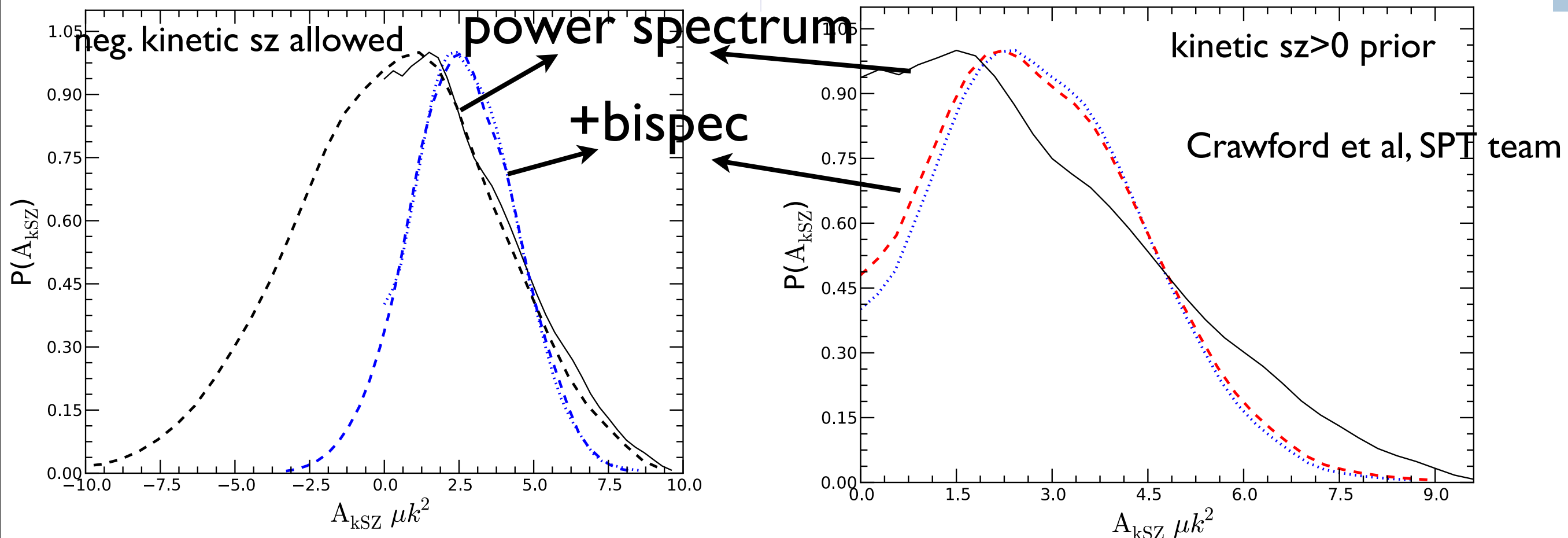
=> combine the bispectrum and power spectrum measurements to individually measure tSZ and kSZ amplitude

=> tSZ amplitude $2.96 \mu k^2 \pm 0.642$ (nominal) ± 0.768 (extreme)

=> improves the tSZ amplitude by factor ~ 2 compared to the power spectrum only case (uncertainty $\sim 1.05 \mu k^2$)



Current SPT measurements: kinetic SZ



95% upper limit for A_{kSZ}

- => power, neg. kinetic sz allowed = $6.50 \mu k^2$
- => power+bispec, neg. kinetic sz allowed = $5.57 \mu k^2$
- => power, $A_{kSZ} > 0$ prior = $6.68 \mu k^2$
- => power+bispec, $A_{kSZ} > 0$ prior(nominal) = $5.59 \mu k^2$
- => power+bispec, $A_{kSZ} > 0$ prior(extreme) = $5.73 \mu k^2$
- => detect $A_{kSZ} = 2.9 \pm 1.5 \mu k^2$

=> Adding bispectrum improves upper limit of the kSZ amplitude by $\sim 15\%$ and we start to see the peak ($A_{kSZ} > 0$ prior case)



Science Cases:

- SZ bispectrum is a new (and powerful) technique to measure the thermal SZ amplitude
- More robust than power spectrum-> signal comes from massive clusters, theoretical uncertainty less compared to power spectrum, kSZ bispectrum is approximately 0, point sources bispectrum is comparable
- A combination of bispectrum+power spectrum measurement can **improve the measurement of thermal and kinetic SZ amplitude individually.**
- A measurement of the kSZ amplitude can provide useful **insight to the reionization epoch.**

