ISOCURVATURE PERTURBATIONS AND REHEATING

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Based on work with Ian Huston, arXiv:1111.6919 and arXiv:1302.4298



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INFLATION: À VERY BRIEF REVIEW

- 'Invented' to solve problems of the big bang:
 - Horizon problem
 - Monopole problem
 - Flatness problem
- Period of rapid accelerated expansion in the early universe

Inflation
$$\iff \ddot{a} > 0 \iff P < -\frac{1}{3}\rho$$

Accelerated expansion driven by scalar field

$$\mathcal{L}_{\varphi} = \frac{1}{2} g^{\mu\nu} \partial_{\mu} \varphi \partial_{\nu} \varphi - V(\varphi)$$

• Homogeneous field:

$$P_{\varphi} = \frac{1}{2}\dot{\varphi}^2 - V(\varphi)$$
$$\rho_{\varphi} = \frac{1}{2}\dot{\varphi}^2 + V(\varphi)$$

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Gives primordial spectrum of density perturbations!

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'simplest' models, specify potential

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increasing complexity

$$\mathcal{L}_{\varphi} = \frac{1}{2} g^{\mu\nu} \partial_{\mu} \varphi \partial_{\nu} \varphi - V(\varphi)$$

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etc...

Zero Parameter Models

3.1 Higgs Inflation (HI)

One Parameter Models

- 4.1 Radiatively Corrected Higgs Inflation (RCHI)
- 4.2 Large Field Inflation (LFI)
- 4.3 Mixed Large Field Inflation (MLFI)
- 4.4 Radiatively Corrected Massive Inflation (RCMI)
- 4.5 Radiatively Corrected Quartic Inflation (RCQI)
- 4.6 Natural Inflation (NI)
- 4.7 Exponential SUSY Inflation (ESI)
- 4.8 Power Law Inflation (PLI)
- 4.9 Kähler Moduli Inflation I (KMII)
- 4.10 Horizon Flow Inflation at first order (HF1I)
- 4.11 Colemann-Weinberg Inflation (CWI)
- 4.12 Loop Inflation (LI)
- 4.13 $(R + R^{2p})$ Inflation (RpI)
- 4.14 Double-Well Inflation (DWI)
- 4.15 Mutated Hilltop Inflation (MHI)
- 4.16 Radion Gauge Inflation (RGI)
- 4.17 MSSM Inflation (MSSMI)
- 4.18 Renormalizable Inflection Point Inflation (RIPI)
- 4.19 Arctan Inflation (AI)
- 4.20 Constant n_s A Inflation (CNAI)
- 4.21 Constant n_s B Inflation (CNBI)
- 4.22 Open String Tachyonic Inflation (OSTI)
- 4.23 Witten-O'Raifeartaigh Inflation (WRI)

Two Parameters Models

- 5.1 Small Field Inflation (SFI)
- 5.2 Intermediate Inflation (II)
- 5.3 Kähler Moduli Inflation II (KMIII)
- 5.4 Logamediate Inflation (LMI)
- 5.5 Twisted Inflation (TWI)
- 5.6 Generalized MSSM Inflation (GMSSMI)

Encyclopædia Inflationaris Page 1/326!!

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- 5.7 Generalized Renormalizable Point Inflation (GRIPI)
- 5.8 Brane SUSY breaking Inflation (BSUSYBI)
- 5.9 Tip Inflation (TI)

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5.10 β exponential inflation (BEI)

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- 5.11 Pseudo Natural Inflation (PSNI)
- 5.12 Non Canonical Kähler Inflation (NCKI)
- 5.13 Constant Spectrum Inflation (CSI)
- 5.14 Orientifold Inflation (OI)
- 5.15 Constant n_s C Inflation (CNCI)
- 5.16 Supergravity Brane Inflation (SBI)
- 5.17 Spontaneous Symmetry Breaking Inflation (SSBI)
- 5.18 Inverse Monomial Inflation (IMI)

Three parameters Models

- 6.1 Running-mass Inflation (RMI)
- 6.2 Valley Hybrid Inflation (VHI)
- 6.3 Dynamical Supersymmetric Inflation (DSI)
- 6.4 Generalized Mixed Inflation (GMLFI)
- 6.5 Logarithmic Potential Inflation (LPI)
- 6.6 Constant n_s D Inflation (CNDI)

Only single field models!

ISOCURVATURE/ENTROPY/ NON-ADIABATIC PRESSURE PERTURBATIONS

ISOCURVATURE/ENTROPY/NON-ADIABATIC PRESSURE PERTURBATIONS

• Adiabatic system $\frac{\delta P}{\dot{P}} = \frac{\delta \rho}{\dot{\rho}}$

 Breaking requirement of adiabaticity allows the pressure perturbation to be expanded as

• The non-adiabatic pressure perturbation can then be defined as $\delta P_{nad} = \delta P - c_s^2 \delta \rho$

ENTROPY PERTURBATIONS

- Thus require a fluid with equation of state depending on two independent variables, i.e.
 P ≡ P(ρ, S)
- Inflationary fields can be written as fluids:
 - single field has barotropic equation of state and so, on large scales, no isocurvature
 - multiple fields, entropy perturbations can be non-zero, and may be important

WHY ISOCURVATURE? 1: VORTICITY

- In classical fluid dynamics, vorticity: $\vec{\omega} = \vec{\nabla} \times \vec{u}$
- Evolves according to

$$\frac{\partial \boldsymbol{\omega}}{\partial t} = \nabla \times (\boldsymbol{v} \times \boldsymbol{\omega}) + \frac{1}{\rho^2} \nabla \rho \times \nabla P$$

- 'Source term' zero if $\vec{\nabla}P$ and $\vec{\nabla}\rho$ are parallel
 - i.e. barotropic fluid, no source term
 - including entropy allows for a source term

Crocco (1937)

VORTICITY IN COSMOLOGY

- Relativistic generalisation, consider vorticity tensor in a perturbed FLRW universe
- Linear perturbations have no source term and decay with expansion

$$\omega_{1ij}' - 3\mathcal{H}c_{\rm s}^2 w_{1ij} = 0$$

Kodama & Sasaki (1984)

VORTICITY IN COSMOLOGY

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Kodama & Sasaki (1984)

 Beyond linear theory source term depends upon the entropy perturbations

$$\omega_{2ij}' - 3\mathcal{H}c_{\rm s}^2\omega_{2ij} = \frac{2a}{\rho_0 + P_0} \left\{ 3\mathcal{H}V_{1[i}\delta P_{{\rm nad}1,j]} + \frac{\delta\rho_{1,[j}\delta P_{{\rm nad}1,i]}}{\rho_0 + P_0} \right\}$$

AJC, Malik & Matravers (2009)

WHY ISOCURVATURE? 2: CMB?



ISOCURVATURE PERTURBATIONS

 We consider multi-field inflationary models with Lagrangian

$$\mathcal{L}(\varphi,\chi) = \frac{1}{2}(\dot{\varphi}^2 + \dot{\chi}^2) + U(\varphi,\chi)$$

with double quartic, quadratic and product exponential potentials

- Study isocurvature perturbations
 - after inflation

- after period of perturbative reheating

MULTI-FIELD INFLATION

• Double quadratic

$$V(\varphi,\chi) = \frac{1}{2}m_{\varphi}^2\varphi^2 + \frac{1}{2}m_{\chi}^2\chi^2$$

• Double quartic $V(\varphi, \chi) = \Lambda^4 \left[\left(1 - \frac{\chi_0^2}{v^2} \right)^2 + \frac{\varphi^2}{\mu^2} + \frac{2\varphi^2\chi^2}{\varphi_c^2v^2} \right]_{\text{Avgoustidis et. al. (2011)}}$ • Product exponential

$$V(\varphi,\chi) = V_0 \varphi^2 e^{-\lambda \chi^2}$$

Byrnes, Choi & Hall (2008)

• Solve background field dynamics and linear perturbations numerically $\ddot{\varphi} + 3H\dot{\varphi} + V_{,\varphi} = 0$

Huston & AJC, PRD 85, 063507 (2012)

http://pyflation.ianhuston.net/

Compute non-adiabatic pressure perturbation

$$\delta P_{\rm nad} = \delta P - c_{\rm s}^2 \delta \rho$$

by defining the pressure and energy density perturbation of the scalar fields as

$$\delta P = \sum_{\alpha} \left(\dot{\varphi_{\alpha}} \delta \dot{\varphi_{\alpha}} - \dot{\varphi_{\alpha}}^2 \phi - V_{,\alpha} \delta \varphi_{\alpha} \right)$$
$$\delta \rho = \sum_{\alpha} \left(\dot{\varphi_{\alpha}} \delta \dot{\varphi_{\alpha}} - \dot{\varphi_{\alpha}}^2 \phi + V_{,\alpha} \delta \varphi_{\alpha} \right)$$

DOUBLE QUADRATIC



Huston & AJC (2012)

DOUBLE QUARTIC



PRODUCT EXPONENTIAL



REHEATING

- How do the scalar fields driving inflation decay into the standard model?
- Include an additional friction term in the scalar field equation $\ddot{\varphi} + (3H + \frac{1}{2}\Gamma)\dot{\varphi} + U_{,\varphi} = 0$
- Usually, one considers the decay of the fields into radiation
- In order to allow isocurvature to survive, consider decay into **both** radiation and matter, subject to constraint $\frac{\Gamma_{m}}{\Gamma_{m}} < 10^{-6}$

- Evolution equations from energy-momentum conservation $\nabla_{\mu}T^{\mu}_{(\alpha)\nu} = Q^{(\alpha)}_{\nu}$
- Fluid equations $\dot{\rho}_{\gamma} + 4H\rho_{\gamma} = \frac{1}{2} \left(\Gamma^{\varphi}_{\gamma} \dot{\varphi}^2 + \Gamma^{\chi}_{\gamma} \dot{\chi}^2 \right)$ $\dot{\rho}_{m} + 3H\rho_{m} = \frac{1}{2} \left(\Gamma^{\varphi}_{m} \dot{\varphi}^2 + \Gamma^{\chi}_{m} \dot{\chi}^2 \right)$

Field equations

$$\ddot{\varphi} + \left[3H + \frac{1}{2}\left(\Gamma_{\gamma}^{\varphi} + \Gamma_{\rm m}^{\varphi}\right)\right]\dot{\varphi} + U_{,\varphi} = 0$$

And similarly (though more complicated) for perturbations

- Solve the evolution equations for the scalar field and matter/radiation fluids
- Again, non-adiabatic pressure perturbation is computed, defined as $\delta P_{nad} = \delta P c_s^2 \delta \rho$
- Now, contribution from the scalar fields, but also from the fluids, i.e.

$$\delta P_{\rm nad} = \frac{1}{3} \delta \rho_{\gamma} \left[1 - \frac{\dot{\rho_{\gamma}}}{\dot{\rho_{\rm m}} + \dot{\rho_{\gamma}}} \right] - \frac{1}{3} \frac{\dot{\rho_{\gamma}} \delta \rho_{\rm m}}{\dot{\rho_{\rm m}} + \dot{\rho_{\gamma}}}$$

Huston & AJC, arXiv:1302.4298

DOUBLE QUADRATIC INFLATION





DOUBLE QUARTIC INFLATION





PRODUCT EXPONENTIAL





SUMMARY

- Isocurvature perturbations common in multifield models; can source source vorticity?
- Reheating should be taken into account when isocurvature perturbations are present
- Non-adiabatic pressure perturbation subdominant in all three models studied

Huston & AJC, PRD 85, 063507 (2012) Huston & AJC, 1302.4298 (2013)

• BUT parameter choice conservative: could we saturate the bounds to source isocurvature?